## Fundamental Interactions

Problem 8: Two-particle decay For the decay $1 \rightarrow 2+3$, where particle 1 is assumed to be at rest, the decay rate is given by

$$
\begin{equation*}
\Gamma=\frac{S}{32 \pi^{2} m_{1}} \int|\mathcal{M}|^{2} \frac{\delta^{4}\left(p_{1}-p_{2}-p_{3}\right)}{\sqrt{\vec{p}_{2}^{2}+m_{2}^{2}} \sqrt{\vec{p}_{3}^{2}+m_{3}^{2}}} d^{3} \vec{p}_{2} d^{3} \vec{p}_{3} \tag{1}
\end{equation*}
$$

where $m_{i}$ is the mass of the $i$ th particle, $p_{i}$ is its four-momentum, and $\vec{p}_{i}$ is its spatial momentum. $S$ is a statistical factor that corrects double-counting when there are identical particles in the final state: i.e. if particle 2 and 3 are identical then $S=1 / 2$ !. The dynamics of the decay process is contained in the amplitude $\mathcal{M}\left(p_{1}, p_{2}, p_{3}\right)$, which we assume to be averaged over spin degrees of freedom.
(i) Express the 4-dimensional delta function into the temporal and the 3-dimensional spatial delta function. Employing that particle 1 is at rest, perform the integral over $\vec{p}_{3}$.
(ii) The amplitude originally depended on all three four-momenta. However, $p_{1}$ is constant for the integration, and $p_{3}$ has been taken care of in the previous step. Moreover, $\mathcal{M}$ must be a Lorentz scalar, such that it can only depend on $\vec{p}_{2}^{2}$.
For the remaining integral change to spherical coordinates $(r, \theta, \phi)$ for $\vec{p}_{2}$ and perform the angular integration.
(iii) Simplify the argument of the remaining 1-dimensional delta function by a change of variable to

$$
\begin{equation*}
u=\sqrt{r^{2}+m_{2}^{2}}+\sqrt{r^{2}+m_{3}^{2}} . \tag{2}
\end{equation*}
$$

Now, evaluate the final integral and verify that

$$
\begin{equation*}
\Gamma=\frac{S\left|\vec{p}_{2}\right|}{8 \pi \mathrm{~m}_{1}^{2}}\left|\mathcal{M}\left(\overrightarrow{\mathrm{p}}_{2}^{2}\right)\right|^{2}, \tag{3}
\end{equation*}
$$

where the formula is evaluated at the particular value

$$
\begin{equation*}
\left|\vec{p}_{2}\right|=\frac{1}{2 m_{1}} \sqrt{m_{1}^{4}+m_{2}^{4}+m_{3}^{4}-2 m_{1}^{2} m_{2}^{2}-2 m_{1}^{2} m_{3}^{2}-2 m_{2}^{2} m_{3}^{2}} \tag{4}
\end{equation*}
$$

determined from the conservation laws.
Without ever knowing the functional form of $\mathcal{M}$ we have been able to carry out the integrals for the 2-body decay. Formula (3) is sometimes referred to as golden rule for a 2-body decay.

Problem 9: Z width We recall that the Standard Model Lagrangian contains the following interaction vertex

$$
\begin{equation*}
\frac{-i g_{2}}{2 \cos \theta_{w}} \gamma^{\mu}\left(c_{V}(f)-c_{A}(f) \gamma^{5}\right) \tag{5}
\end{equation*}
$$

between the $Z$ boson and a fermion anti-fermion pair $f \bar{f} . c_{V}\left(c_{A}\right)$ denotes the vector (axial-vector) coupling of the fermion to the Z-boson and are given by

$$
\begin{align*}
& c_{V}(f)=T_{3}(f)-2 Q(f) \sin ^{2} \theta_{W}  \tag{6a}\\
& c_{A}(f)=T_{3}(f) . \tag{6b}
\end{align*}
$$

| Standard Model fermions |  | Q | $\mathrm{T}_{3}$ | $\mathrm{~m}_{\mathrm{f}}$ (approximate) |
| :---: | :---: | :---: | :---: | :---: |
| leptons | $e^{-}, \mu^{-}, \tau^{-}$ | -1 | $-\frac{1}{2}$ | $511 \mathrm{keV}, 105 \mathrm{MeV}, 1776 \mathrm{MeV}$ |
| neutrinos | $v_{e}, \nu_{\mu}, v_{\tau}$ | 0 | $\frac{1}{2}$ | $<2 \mathrm{eV}$ |
| up-type quarks | $\mathrm{u}, \mathrm{c}, \mathrm{t}$ | $\frac{2}{3}$ | $\frac{1}{2}$ | $2,3 \mathrm{MeV}, 1,275 \mathrm{GeV}, 173,5 \mathrm{GeV}$ |
| down-type quarks | $\mathrm{d}, \mathrm{s}, \mathrm{b}$ | $-\frac{1}{3}$ | $-\frac{1}{2}$ | $4,8 \mathrm{MeV}, 95 \mathrm{MeV}, 4,2 \mathrm{GeV}$ |

Table 1: Particle properties, cf. Particle data group http://pdg.lbl.gov.

Here, $T_{3}$ denotes the third component of the weak isospin and $Q$ the electric charge. These properties are provided in Tab. 1 for the Standard Model fermions. The amplitude $\mathcal{M}_{Z \rightarrow f \bar{f}}$ is computed by multiplying the fermion (anti-fermion) wave-function from the right(left), as well as the polarisation vector for the $Z$-boson, and averaging over the spin and color space. In the approximation that $m_{f} \ll m_{Z}$ ( $m_{Z}$ denotes the Z-boson mass), one obtains

$$
\begin{equation*}
\left|\mathcal{M}_{Z \rightarrow \bar{f}}\right|^{2}=\frac{N_{c}}{3} \frac{g_{2}^{2}}{\cos ^{2} \theta_{W}} m_{Z}^{2}\left(c_{V}^{2}(f)+c_{\mathcal{A}}^{2}(f)\right) \tag{7}
\end{equation*}
$$

and $N_{c}$ is the number of different color states the fermion $f$ can occupy.
i) What are the fermion anti-fermion pairs relevant for the decay $Z \rightarrow f \bar{f}$ ? You should identify four distinct cases.
ii) Employing (3) in the limit $\mathfrak{m}_{f} \ll m_{Z}$ and (7), verify that the decay width $\Gamma_{f \bar{f}}$ of the $Z$-boson into a fermion anti-fermion pair is given by

$$
\begin{equation*}
\Gamma_{\mathrm{f} \overline{\mathrm{f}}}=\frac{\mathrm{N}_{\mathrm{c}}}{48 \pi} \frac{\mathrm{~g}_{2}^{2}}{\cos ^{2} \theta_{W}} m_{\mathrm{Z}}\left(\mathrm{c}_{\mathrm{V}}^{2}+\mathrm{c}_{\mathrm{A}}^{2}\right) \tag{8}
\end{equation*}
$$

iii) Evaluate $c_{V}, c_{A}$ for the four cases and provide the explicit formula for $\Gamma_{\bar{f} \bar{f}}$.
iv) Using $\sin ^{2} \theta_{w}=0,23, g_{2}=\frac{e}{\sin \theta_{w}}$ (or more conveniently $g_{2}^{2} \approx \frac{4 \pi}{30}$ ), and $m_{Z}=91,187 \mathrm{GeV}$, compute the numerical values for the partial widths $\Gamma_{f \bar{f}}$ as well as the total Z decay width

$$
\begin{equation*}
\Gamma_{Z}=\sum_{f} \Gamma_{f \bar{f}} . \tag{9}
\end{equation*}
$$

