

Exercises V

November 19th

E5.1 The Born Approximation

The Born approximation for (elastic) potential scattering in first order reads

$$\frac{d\sigma}{d\Omega} = \frac{2\pi m^2}{\hbar^4} \left| \int d^3r \frac{1}{(2\pi)^{\frac{3}{2}}} V(\vec{r}) e^{i\vec{r}\cdot(\vec{k}_i - \vec{k}_f)} \right|^2, \quad \text{where } |\vec{k}_f| = |\vec{k}_i| = k. \quad (*)$$

Briefly stated, the differential cross section is given by the Fourier transform of the potential, taken at the value of the momentum transfer, i.e.

$$\frac{d\sigma}{d\Omega} = \frac{2\pi m^2}{\hbar^4} |(\mathcal{F}V)(\vec{q})|^2, \quad \vec{q} = \vec{k}_i - \vec{k}_f.$$

Anticipating the lecture, deduce (*) from the Golden Rule with your tutor's help.

- (1) Show for the momentum transfer that $|\vec{q}| = 2k \sin \frac{\vartheta}{2}$. ϑ is the scattering angle.
- (2) Perform the angular integration in (*) for a radially symmetric potential:

$$\frac{d\sigma}{d\Omega} = \frac{4m^2}{\hbar^4} \frac{1}{|\vec{q}|^2} \left| \int_0^\infty dr r V(r) \sin(|\vec{q}|r) \right|^2.$$

- (3) Consider a Yukawa potential $V(r) = \frac{\kappa}{r} e^{-\frac{r}{r_0}}$. Zeige mit (2), daß

$$\frac{d\sigma}{d\Omega} = \left(\frac{\kappa}{4E \sin^2 \frac{\vartheta}{2} + \frac{\hbar^2}{2mr_0^2}} \right)^2.$$

- (4) Which famous formula does one obtain in (3) taking the limit of the Coulomb potential? Does \hbar still occur? Discuss the case of forward scattering.
- (5) Show for a Gaussian potential $V(r) = V_0 \exp(-\frac{r^2}{2r_0^2})$ that

$$\frac{d\sigma}{d\Omega} = \frac{2\pi m^2 r_0^6 V_0^2}{\hbar^4} e^{-4k^2 r_0^2 \sin^2 \frac{\vartheta}{2}}.$$

How does this cross section behave for large scattering angles?

- (6) Show for a spherically symmetric potential well with depth V_0 and radius r_0 :

$$\frac{d\sigma}{d\Omega} = \frac{m^2 V_0}{\pi \hbar^4} \frac{1}{|\vec{q}|^4} \left(r_0 \cos(|\vec{q}|r_0) - \frac{\sin(|\vec{q}|r_0)}{|\vec{q}|} \right)^2.$$

How does this cross section behave for large scattering energies?

- (7) Discuss $\frac{d\sigma}{d\Omega}$ in (3), (5) and (6) for small energies ($kr_0 \ll 1$). Can one distinguish these potentials by scattering slow projectiles? What do they have in common?
- (8) Let $V(\vec{r})$ be the Coulomb potential of a charge distribution ϱ , i.e.

$$V(\vec{r}) = Q \int d^3r' \frac{\varrho(\vec{r}')}{|\vec{r} - \vec{r}'|}.$$

Show

$$\frac{d\sigma}{d\Omega} = \frac{32\pi^3 m^2 Q^2}{\hbar^4} \frac{1}{|\vec{q}|^4} |(\mathcal{F}\varrho)(\vec{q})|^2.$$

$|(\mathcal{F}\varrho)(\vec{q})|^2$ is called *form factor*. What does one obtain for a pointlike charge?

Homework V

Return: November 26th

H5.1 *Pion–proton scattering*

- (1) Compute $\frac{d\sigma}{d\Omega}$ for the superposition of a Yukawa and a Coulomb potential:

$$\frac{d\sigma}{d\Omega} = \left(\frac{\kappa}{4E \sin^2 \frac{\vartheta}{2} + \frac{\hbar^2}{2mr_0^2}} + \frac{q_1 q_2}{4E \sin^2 \frac{\vartheta}{2}} \right)^2.$$

- (2) What will happen if the detector is placed in forward direction ($\vartheta = 0$)?
 (3) Compare the three contributions of the Coulomb potential, the Yukawa potential and their interference in $\frac{d\sigma}{d\Omega}$ depending on E for $\sin \frac{\vartheta}{2} > 0.1$, i.e. away from the forward direction. Let $\kappa = 0.07 \hbar c$, $r_0 = 1.4 \text{ fm}$, $mc^2 = 140 \text{ MeV}$, $q_1 = q_2 = e$.
 (8 points)

H5.2 *Scattering of an electron by a hydrogen atom*

Let the atom be in the ground state. At first, consider elastic scattering.

- (1) Compute the total charge density for the hydrogen atom to be

$$\rho(\vec{r}) = e\delta^{(3)}(\vec{r}) - \frac{e}{\pi r_0^3} e^{-\frac{2r}{r_0}}.$$

- (2) Using E5.1.8, compute the differential cross section to be

$$\frac{d\sigma}{d\Omega} = \frac{4}{r_0^2} \frac{1}{|\vec{q}|^4} \left(1 - \frac{16}{(4 + |\vec{q}|^2 r_0^2)^2} \right)^2.$$

Now take the scattering to be inelastic, i.e. the atom can be excited to $|nlm\rangle$.

- (3) Show by transferring the derivation of the formula for the differential cross section for elastic scattering to the inelastic case that

$$\begin{aligned} \frac{d\sigma}{d\Omega} = & \frac{2\pi m}{\hbar^2 |\vec{k}_i|} \int dk k^2 \delta \left(\frac{\hbar^2 |\vec{k}|^2}{2m} + E_{n\ell} - \frac{\hbar^2 |\vec{k}_i|^2}{2m} - E_{1s} \right) \\ & \times \left| \iint \frac{1}{(2\pi)^{\frac{3}{2}}} \psi_{nlm}^*(\vec{r}') \psi_{1s}(\vec{r}') \left(-\frac{e^2}{r} + \frac{e^2}{|\vec{r} - \vec{r}'|} \right) e^{i\vec{r} \cdot (\vec{k}_i - \vec{k})} d^3 r d^3 r' \right|^2. \end{aligned}$$

Here, k_i denotes the wave number vector of the incoming particle.

- (4) Why does the potential of the nucleus not contribute to the inelastic scattering?
 (5) Show in (3) by applying the convolution theorem that

$$\frac{d\sigma}{d\Omega} = \frac{32\pi^3 m^2 e^4}{\hbar^4} \frac{|\vec{k}_f|}{|\vec{k}_i|} \frac{1}{|\vec{q}|^4} |(\mathcal{F}(\psi_{nlm}^* \psi_{1s}))(\vec{q})|^2 \quad \text{with} \quad |\vec{k}_f|^2 = \frac{2m}{\hbar^2} \left(E_{1s} - E_{n\ell} + \frac{\hbar^2 |\vec{k}_i|^2}{2m} \right),$$

$$|\vec{q}|^2 = 2 \left(|\vec{k}_i|^2 + \frac{(E_{1s} - E_{n\ell})m}{\hbar^2} - \frac{|\vec{k}_i|}{\hbar} \sqrt{m \left(2(E_{1s} - E_{n\ell}) + \frac{\hbar^2 |\vec{k}_i|^2}{m} \right) \cos \vartheta} \right).$$

- (6) Compute $\frac{d\sigma}{d\Omega}$ explicitly for the case that the atom is excited to the $2s$ state.

(22 points)