

Lecture 17

Spherically symmetric metrics. Calculation of Riemann-, Ricci-, and Einstein-Tensors using the Cartan Structure equations.

Eventually we seek solutions to Einstein's equations for the idealised case of a spherically symmetric star; inside ($T_{\mu\nu} \neq 0$) and outside ($T_{\mu\nu} = 0$) the star. But first we need to define "spherical symmetry".

Definition: Let (M, g) be a 4-dimensional spacetime and

$$\left. \begin{aligned} \phi &: SO(3) \rightarrow \text{Diff}(M) \\ h &\mapsto \phi_h \end{aligned} \right\} \quad (17.1)$$

be an isometric action of $SO(3)$; that is

$$\left. \begin{aligned} \phi_e &= \text{id}_M \\ \phi_{h'} \circ \phi_h &= \phi_{h' \cdot h} \end{aligned} \right\} \quad \text{"action"} \quad (17.2)$$

with

$$\left. \begin{aligned} \phi_h^* g &= g \end{aligned} \right\} \quad \text{"isometric"} \quad (17.3)$$

$$\forall h \in SO(3)$$

We further require the orbits of this action to be spacelike S^2 (2-spheres); that is

$$\text{Orb}_p(SO(3)) = \bigcup_{h \in SO(3)} \phi_h(p) \quad (17.4)$$

is an embedded spacelike topological 2-sphere whose geometry is necessarily "round", i.e. isometric to that of a standard 2-sphere in \mathbb{R}^3 of some radius. This is because of our previous requirement (17-3) that $SO(3)$ acts by isometries.

A spacetime admitting an isometric action of $SO(3)$ with spacelike S^2 -orbits is called spherically symmetric.

Theorem: A spherically symmetric spacetime admits local charts in which

$$g = \left. \begin{aligned} &e^{2a} \dot{c}^2 dt \otimes dt - e^{2b} dr \otimes dr \\ &- r^2 (d\theta \otimes d\theta + \sin^2(\theta) d\varphi \otimes d\varphi) \end{aligned} \right\} (17.5)$$

where a and b depend on (t, r) but not on θ, φ .

The $SO(3)$ orbits are given by the 2-spheres of constant t and r , which are isometric to the 2-spheres of metric radius r in \mathbb{R}^3 . Their area in (M, g) is

$$A(r) = 4\pi r^2 \quad (17.6)$$

just like for a sphere in \mathbb{R}^3 of euclidean radius r . It is for this reason that our coordinate r is called "area radius".

It may be geometrically defined by

$$r : M \rightarrow \mathbb{R}_+$$

$$p \mapsto r(p) := \left(\frac{A(p)}{4\pi} \right)^{1/2} \quad (17.7)$$

where

$$A(p) = \text{Area}(\text{Orb}_p(SO(3))). \quad (17.8)$$

Now, for the theorem above to be valid we must assume that this function is a coordinate, which is only true iff

$$dr \neq 0 \quad (17.9)$$

This condition must be added to the list of hypotheses for our theorem.

We shall consider the general case in an exercise.

Given a metric of the form (17.5) we can now write it in terms of orthonormal co-basis:

$$g = \Theta^0 \otimes \Theta^0 - \sum_{a=1}^3 \Theta^a \otimes \Theta^a \quad (17.10)$$

$$\Theta^0 = e^a c dt = e^a dx^0$$

$$\Theta^1 = e^b dr$$

$$\Theta^2 = r d\theta$$

$$\Theta^3 = r \sin\theta d\varphi$$

(17.11)

where $a = a(t, r)$, $b = b(t, r)$.

We shall write

$$\dot{a} := \frac{\partial a}{\partial x^0}, \quad a' := \frac{\partial a}{\partial r}$$

$$\dot{b} := \frac{\partial b}{\partial x^0}, \quad b' := \frac{\partial b}{\partial r}$$

(17.12)

We now employ the Cartan Structure Equations to calculate all connection coefficients and all Curvature components with respect to this orthonormal frame.

The 1st Structure equation is ($T=0$)

$$d\Theta^\alpha + \omega^\alpha{}_\beta \wedge \Theta^\beta = 0 \quad (17.13)$$

Where for $\omega_{\alpha\beta} = \eta_{\alpha\gamma} \omega^\gamma{}_\beta$ we have

$$\omega_{\alpha\beta} = -\omega_{\beta\alpha} \quad (17.14)$$

which expresses the metricity of the connection.

We proceed by evaluating (17.13) for $\alpha = 0, 1, 2, 3$. Using (17.11):

$\alpha = 0$)

$$\begin{aligned} d\Theta^0 &= e^a (\dot{a} dx^0 + a' dt) \wedge dx^0 \\ &= -a' e^{-b} \Theta^0 \wedge \Theta^1 \\ &= -\omega^0{}_a \wedge \Theta^a \quad (\text{note } \omega^0{}_0 = 0) \end{aligned}$$

$$\Rightarrow \omega^0{}_1 = a' e^{-b} \Theta^0 + \text{Terms} \sim \Theta^1 \quad (17.15a)$$

$$\omega^0{}_2 = \text{Terms} \sim \Theta^2 \quad (17.15b)$$

$$\omega^0{}_3 = \text{Terms} \sim \Theta^3 \quad (17.15c)$$

$d=1)$

$$\begin{aligned}
 d\Theta^1 &= e^b (\dot{b} dx^0 + b' dt) \wedge dr \\
 &= -\dot{b} e^{-a} \Theta^1 \wedge \Theta^0 \\
 &= -\omega^1_0 \wedge \Theta^0 - \omega^1_2 \wedge \Theta^2 - \omega^1_3 \wedge \Theta^3
 \end{aligned}$$

$$\Rightarrow \omega^1_0 = \dot{b} e^{-a} \Theta^1 + \text{Terms} \sim \Theta^0 \quad (17.16a)$$

$$\omega^1_2 = \text{Terms} \sim \Theta^2 \quad (17.16b)$$

$$\omega^1_3 = \text{Terms} \sim \Theta^3 \quad (17.16c)$$

 $d=2)$

$$\begin{aligned}
 d\Theta^2 &= dt \wedge d\Theta = \frac{1}{r} e^{-b} \Theta^1 \wedge \Theta^2 \\
 &= -\omega^2_0 \wedge \Theta^0 - \omega^2_1 \wedge \Theta^1 - \omega^2_3 \wedge \Theta^3
 \end{aligned}$$

$$\Rightarrow \omega^2_0 = \text{Terms} \sim \Theta^0 \quad (17.17a)$$

$$\omega^2_1 = \frac{1}{r} e^{-b} \Theta^2 + \text{Terms} \sim \Theta^1 \quad (17.17b)$$

$$\omega^2_3 = \text{Terms} \sim \Theta^3 \quad (17.17c)$$

 $d=3)$

$$\begin{aligned}
 d\Theta^3 &= \sin\theta dr \wedge d\varphi + r \cos\theta d\theta \wedge d\varphi \\
 &= \frac{1}{r} e^{-b} \Theta^1 \wedge \Theta^3 + \frac{1}{r} \cot\theta \Theta^2 \wedge \Theta^3 \\
 &= -\omega^3_0 \wedge \Theta^0 - \omega^3_1 \wedge \Theta^1 - \omega^3_2 \wedge \Theta^2
 \end{aligned}$$

$$\Rightarrow \omega^3_0 = \text{Terms} \sim \Theta^0 \quad (17.18a)$$

$$\omega^3_1 = \frac{1}{r} e^{-b} \Theta^3 + \text{Terms} \sim \Theta^1 \quad (17.18b)$$

$$\omega^3_2 = \frac{1}{r} \cot \Theta \Theta^3 + \text{Terms} \sim \Theta^2 \quad (17.18c)$$

Using

$$\omega^0_a = \omega^a_0$$

and $\omega^a_b = -\omega^b_a$

We get from

(17.15a) and (17.16a):

$$\omega^0_1 = \omega^1_0 = a^1 e^{-b} \Theta^0 + b^1 e^{-a} \Theta^1 \quad (17.19a)$$

(17.15b) and (17.17a):

$$\omega^0_2 = \omega^2_0 = 0 \quad (17.19b)$$

(17.15c) and (17.18a):

$$\omega^0_3 = \omega^3_0 = 0 \quad (17.19c)$$

(17.16b) and (17.17b):

$$\omega^1_2 = -\omega^2_1 = -\frac{1}{r} e^{-b} \Theta^2 \quad (17.19d)$$

(17.16c) and (17.18b):

$$\omega^1_3 = -\omega^3_1 = -\frac{1}{r} e^{-b} \Theta^3 \quad (17.19e)$$

(17.17c) and (17.18c)

$$\omega^2_3 = -\omega^3_2 = -\frac{1}{r} \cot \theta \theta^3 \quad (17.19e)$$

To summarise we get for the connection coefficients

$$\omega^0_1 = \omega^1_0 = a' e^{-b} \theta^0 + b' e^{-a} \theta^1 \quad (17.20a)$$

$$= a' e^{a-b} dx^0 + b' e^{b-a} dr \quad (17.20b)$$

$$\omega^0_2 = \omega^0_2 = 0 \quad (17.21)$$

$$\omega^0_3 = \omega^3_0 = 0 \quad (17.22)$$

$$\omega^1_2 = -\omega^2_1 = -\frac{1}{r} e^{-b} \theta^2 \quad (17.23a)$$

$$= -e^{-b} d\theta \quad (17.23b)$$

$$\omega^1_3 = -\omega^3_1 = -\frac{1}{r} e^{-b} \theta^3 \quad (17.24a)$$

$$= -e^{-b} \sin \theta d\varphi \quad (17.24b)$$

$$\omega^2_3 = -\omega^3_2 = -\frac{1}{r} \cot \theta \theta^3 \quad (17.25a)$$

$$= -\cos \theta d\varphi \quad (17.25b)$$

The Curvature is now calculated by the 2nd Cartan Structure equation:

$$\Omega^\alpha{}_\beta = d\omega^\alpha{}_\beta + \omega^\alpha{}_\lambda \wedge \omega^\lambda{}_\beta$$

01)

$$\begin{aligned} d\omega^0{}_1 &= (a' e^{a-b})' dr \wedge dx^0 \\ &\quad + (i' e^{b-a})' dx^0 \wedge dr \\ &= e^{-(a+b)} [(i' e^{b-a})' - (a' e^{a-b})'] \theta^0 \wedge \theta^1 \end{aligned}$$

$$\omega^0{}_\lambda \wedge \omega^\lambda{}_1 = \omega^0{}_2 \wedge \omega^2{}_1 + \omega^0{}_3 \wedge \omega^3{}_1 = 0$$

$$\Rightarrow \Omega^0{}_1 = e^{-(a+b)} [(i' e^{b-a})' - (a' e^{a-b})'] \theta^0 \wedge \theta^1 \quad (17.26a)$$

02)

$$d\omega^0{}_2 = 0$$

$$\omega^0{}_\lambda \wedge \omega^\lambda{}_2 = \omega^0{}_1 \wedge \omega^1{}_2 + \omega^0{}_3 \wedge \omega^3{}_2$$

$$= (a' e^{-b} \theta^0 + i' e^{-a} \theta^1) \wedge (-\frac{1}{r} e^{-b}) \theta^2$$

$$= -\frac{a'}{r} e^{-2b} \theta^0 \wedge \theta^2 - \frac{i'}{r} e^{-(a+b)} \theta^1 \wedge \theta^2$$

$$\Rightarrow \Omega^0{}_2 = -\frac{a'}{r} e^{-2b} \theta^0 \wedge \theta^2 - \frac{i'}{r} e^{-(a+b)} \theta^1 \wedge \theta^2 \quad (17.26b)$$

03)

$$d\omega^0_3 = 0$$

$$\omega^0_1 \wedge \omega^1_3 = \omega^0_1 \wedge \omega^1_3$$

$$= (a' e^{-b} \theta^0 + i b' e^{-a} \theta^1) (-\frac{1}{r} e^{-b} \theta^3)$$

$$= -\frac{a'}{r} e^{-2b} \theta^0 \wedge \theta^3 - \frac{i b'}{r} e^{-(a+b)} \theta^1 \wedge \theta^3$$

$$\Rightarrow \Omega^0_3 = -\frac{a'}{r} e^{-2b} \theta^0 \wedge \theta^3 - \frac{i b'}{r} e^{-(a+b)} \theta^1 \wedge \theta^3 \quad (17.26c)$$

12)

$$d\omega^1_2 = e^{-b} (i b' dx^0 + b' dr) \wedge d\theta$$

$$= \frac{i b'}{r} e^{-(a+b)} \theta^0 \wedge \theta^2 + \frac{b'}{r} e^{-2b} \theta^1 \wedge \theta^2$$

$$\omega^1_1 \wedge \omega^1_2 = \omega^1_0 \wedge \cancel{\omega^0_2} + \omega^1_3 \wedge \omega^3_2$$

$$= 0 \quad \text{since } \omega^1_3 \sim \theta^3 \text{ and } \omega^3_2 \sim \theta^3$$

$$\Rightarrow \Omega^1_2 = \frac{i b'}{r} e^{-(a+b)} \theta^0 \wedge \theta^2 + \frac{b'}{r} e^{-2b} \theta^1 \wedge \theta^2 \quad (17.26d)$$

13)

$$d\omega^1_3 = e^{-b} (i b' dx^0 + b' dr) \wedge \sin\theta dy$$

$$- e^{-b} \cos\theta d\theta \wedge dy$$

$$= \frac{b}{r} e^{-(a+b)} \theta^0 \wedge \theta^3 + \frac{b'}{r} e^{-2b} \theta^1 \wedge \theta^3 - \frac{\cot \theta}{r^2} e^{-b} \theta^2 \wedge \theta^3$$

$$\begin{aligned} \omega^1 \wedge \omega^3 &= \omega^1_0 \wedge \omega^3_0 + \omega^1_2 \wedge \omega^3_2 \\ &= \left(-\frac{1}{r} e^{-b} \theta^2\right) \wedge \left(-\frac{1}{r} \cot \theta \theta^3\right) \\ &= + \frac{\cot \theta}{r^2} e^{-b} \theta^2 \wedge \theta^3 \end{aligned}$$

$$\Rightarrow \Omega^1_3 = \frac{b}{r} e^{-(a+b)} \theta^0 \wedge \theta^3 + \frac{b'}{r} e^{-2b} \theta^1 \wedge \theta^3 \quad (17.26e)$$

23)

$$\begin{aligned} d\omega^2_3 &= \sin \theta \, d\theta \wedge d\varphi \\ &= \frac{1}{r^2} \theta^2 \wedge \theta^3 \end{aligned}$$

$$\begin{aligned} \omega^2 \wedge \omega^3 &= \omega^2_0 \wedge \omega^3_0 + \omega^2_1 \wedge \omega^3_1 \\ &= \frac{1}{r} e^{-b} \theta^2 \wedge \left(-\frac{1}{r} e^{-b} \theta^3\right) \\ &= -\frac{1}{r^2} e^{-2b} \theta^2 \wedge \theta^3 \end{aligned}$$

$$\Rightarrow \Omega^2_3 = \frac{1}{r^2} [1 - e^{-2b}] \theta^2 \wedge \theta^3 \quad (17.26f)$$

Flow

$$\Omega^\alpha{}_\beta = \frac{1}{2} R^\alpha{}_{\beta\mu\nu} \Theta^\mu \wedge \Theta^\nu \quad (17.27)$$

With $R^\alpha{}_{\beta\mu\nu} = -R^\alpha{}_{\beta\nu\mu}$, we read off all curvature components in the orthonormal frame Θ^α :

$$\begin{aligned} R^0{}_{101} &= R_{0101} \\ &= e^{-2a} (\ddot{b} + \dot{b}^2 - \dot{b}\dot{a}) - e^{-2b} (a' + a'^2 - a'b') \end{aligned} \quad (17.28a)$$

$$R^0{}_{202} = R_{0202} = -\frac{a'}{\lambda} e^{-2b} \quad (17.28b)$$

$$R^0{}_{212} = R_{0212} = -\frac{b'}{\lambda} e^{-(a+b)} \quad (17.28c)$$

$$R^0{}_{303} = R_{0303} = -\frac{b'}{\lambda} e^{-2b} \quad (17.28d)$$

$$R^0{}_{313} = R_{0313} = -\frac{a'}{\lambda} e^{-(a+b)} \quad (17.28e)$$

$$R^1{}_{212} = -R_{1212} = \frac{b'}{\lambda} e^{-2b} \quad (17.28f)$$

$$R^1{}_{202} = -R_{1202} = \frac{a'}{\lambda} e^{-(a+b)} \quad (17.28g)$$

$$R^1{}_{313} = -R_{1313} = \frac{b'}{\lambda} e^{-2b} \quad (17.28h)$$

$$R^1{}_{303} = -R_{1303} = \frac{a'}{\lambda} e^{-(a+b)} \quad (17.28i)$$

$$R^2{}_{323} = -R_{2323} = \frac{1}{\lambda^2} (1 - e^{-2b}) \quad (17.28j)$$

All other components vanish!

The non vanishing components of the Ricci-Tensor are

$$\begin{aligned}
 R_{00} &= R^\lambda_{\ 0 \times 0} = -R_{1010} - R_{2020} - R_{3030} \\
 &= e^{-2b} (a'' + a'^2 - a'b') - e^{-2a} (\ddot{b} + \dot{b}^2 - \dot{a}\dot{b}) \\
 &\quad + \frac{2a'}{r} e^{-2b} \\
 &= e^{-2b} \left(a'' + a'^2 - a'b' + \frac{2a'}{r} \right) \\
 &\quad - e^{-2a} (\ddot{b} + \dot{b}^2 - \dot{a}\dot{b}) \quad (17.29a)
 \end{aligned}$$

$$\begin{aligned}
 R_{01} &= R^\lambda_{\ 0 \times 1} = -R_{2021} - R_{3031} \\
 &= 2 \frac{\dot{b}}{r} e^{-(a+b)} \quad (17.29b)
 \end{aligned}$$

$$R_{02} = R^\lambda_{\ 0 \times 2} = -R_{1012} - R_{3032} = 0 \quad (17.29c)$$

$$R_{03} = R^\lambda_{\ 0 \times 3} = -R_{1013} - R_{2023} = 0 \quad (17.29d)$$

$$\begin{aligned}
 R_{11} &= R^\lambda_{\ 1 \times 1} = R_{0101} - R_{2121} - R_{3131} \\
 &= e^{-2a} (\ddot{b} + \dot{b}^2 - \dot{a}\dot{b}) - e^{-2b} (a'' + a'^2 - a'b') \\
 &\quad + 2 \frac{b'}{r} e^{-2b} \\
 &= e^{-2a} (\ddot{b} + \dot{b}^2 - \dot{a}\dot{b}) - e^{-2b} \left(a'' + a'^2 - a'b' - 2 \frac{b'}{r} \right) \quad (17.29e)
 \end{aligned}$$

$$R_{12} = R^{\lambda}_{1\lambda 2} = R_{0102} - R_{3132} = 0 \quad (17.29f)$$

$$R_{13} = R^{\lambda}_{1\lambda 3} = R_{0103} - R_{2123} = 0 \quad (17.29g)$$

$$\begin{aligned} R_{22} &= R^{\lambda}_{2\lambda 2} = R_{0202} - R_{1212} - R_{3232} \\ &= -\frac{a'}{r} e^{-2b} + \frac{b'}{r} e^{-2b} + \frac{1}{r^2} (1 - e^{-2b}) \\ &= (b' - a') \frac{1}{r} e^{-2b} + \frac{1}{r^2} (1 - e^{-2b}) \end{aligned} \quad (17.29h)$$

$$R_{23} = R^{\lambda}_{2\lambda 3} = R_{0203} - R_{1213} = 0 \quad (17.29i)$$

$$\begin{aligned} R_{33} &= R^{\lambda}_{3\lambda 3} = R_{0303} - R_{1313} - R_{2323} \\ &= -\frac{a'}{r} e^{-2b} + \frac{b'}{r} e^{-2b} + \frac{1}{r^2} (1 - e^{-2b}) \\ &= (b' - a') \frac{1}{r} e^{-2b} + \frac{1}{r^2} (1 - e^{-2b}) \end{aligned} \quad (17.29j)$$

The Ricci-scalar is

$$\begin{aligned} R &= \eta^{\alpha\beta} R_{\alpha\beta} = R_{00} - R_{11} - R_{22} - R_{33} \\ &= 2 e^{-2b} \left(a'' + a'^2 - a'b' + \frac{a' - b'}{r} \right) \\ &\quad - 2 e^{-2a} \left(\ddot{b} + \dot{b}^2 - \dot{a}\dot{b} \right) \\ &\quad - 2 (b' - a') \frac{1}{r} e^{-2b} - \frac{2}{r^2} (1 - e^{-2b}) \\ &= 2 \left\{ e^{-2b} \left(a'' + a'^2 - a'b' + 2 \frac{a' - b'}{r} + \frac{1}{r^2} \right) \right. \\ &\quad \left. - e^{-2a} \left(\ddot{b} + \dot{b}^2 - \dot{a}\dot{b} \right) - \frac{1}{r^2} \right\} \end{aligned} \quad (17.30)$$

The components of the Einstein-Tensor are easier to determine directly instead of using this long expression for the Ricci Scalar; i.e.

$$\begin{aligned}
 G_{00} &= R_{00} - \frac{1}{2} \eta_{00} (R_{00} - R_{11} - R_{22} - R_{33}) \\
 &= \frac{1}{2} (R_{00} + R_{11} + R_{22} + R_{33}) \\
 &= \frac{1}{2} \left(2 \frac{a'+b'}{r} e^{-2b} + 2 \frac{b'-a'}{r} e^{-2b} + \frac{2}{r^2} (1 - e^{-2b}) \right) \\
 &= \left(\frac{2b'}{r} - \frac{1}{r^2} \right) e^{-2b} + \frac{1}{r^2} \\
 &= \frac{1}{r^2} \left[1 - (r e^{-2b})' \right] \quad (17.31a)
 \end{aligned}$$

$$\begin{aligned}
 G_{11} &= R_{11} - \frac{1}{2} \eta_{11} (R_{00} - R_{11} - R_{22} - R_{33}) \\
 &= \frac{1}{2} (R_{11} + R_{00} - R_{22} - R_{33}) \\
 &= \frac{1}{2} \left(2 \frac{a'+b'}{r} e^{-2b} - 2 \frac{b'-a'}{r} e^{-2b} - \frac{2}{r^2} (1 - e^{-2b}) \right) \\
 &= \left(\frac{2a'}{r} + \frac{1}{r^2} \right) e^{-2b} - \frac{1}{r^2} \quad (17.31b)
 \end{aligned}$$

Looking ahead, we note that

$$G_{00} + G_{11} = (a+b)' \frac{2}{r} e^{-2b}. \quad (17.31c)$$

$$\begin{aligned}
 G_{22} &= R_{22} - \frac{1}{2} \eta_{22} (R_{00} - R_{11} - R_{22} - R_{33}) \\
 &= \frac{1}{2} (R_{00} + \cancel{R_{22}} - R_{11} - \cancel{R_{33}}) \\
 &= \frac{1}{2} (R_{00} - R_{11}) \\
 &= e^{-2b} \left(\dot{a}'' + a'^2 - a'b' + \frac{a' - b'}{r} \right) \\
 &\quad - e^{-2a} (\ddot{b} + \dot{b}^2 - \dot{a}\dot{b}) \quad (17.31d)
 \end{aligned}$$

$$G_{33} = G_{22} \quad (17.31e)$$

$$G_{01} = R_{01} = 2 \frac{\dot{b}}{r} e^{-(a+b)} \quad (17.31f)$$

$$G_{02} = R_{02} = 0$$

$$G_{03} = R_{03} = 0$$

$$G_{12} = R_{12} = 0$$

$$G_{13} = R_{13} = 0$$

$$G_{23} = R_{23} = 0$$

(17.31g)

Hence the only non-diagonal component of the Einstein Tensor is G_{01} (radial energy current density). This is obvious from spherical symmetry, which cannot allow for non-zero vector-field tangent to S^2 .