

Exercises for the lecture on  
**Introduction into General Relativity**  
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**Sheet 10**

**Problem 1**

Consider a 3-parameter family of timelike geodesics  $\gamma(s, \vec{\sigma})$ , where  $s$  is the proper length along the geodesic labelled by  $\vec{\sigma} := (\sigma_1, \sigma_2, \sigma_3)$ . This is sometimes called a (3-dimensional) “geodesic congruence”. We assume that  $\vec{\sigma}$  ranges over some open set  $\mathcal{U} \subset \mathbb{R}^3$  and that the map

$$\gamma : \mathbb{R} \times \mathcal{U} \rightarrow M, \quad (s, \vec{\sigma}) \mapsto \gamma(s, \vec{\sigma}) \quad (1)$$

is a smooth diffeomorphism of  $\mathbb{R} \times \mathcal{U}$  onto the image  $\text{Im}(\gamma) \subseteq M$  of  $\gamma$ .

For any fixed  $\vec{v} \in \mathbb{R}^3$  consider the following vector fields on  $\text{Im}(\gamma)$ :

$$\dot{\gamma} := \frac{\partial \gamma}{\partial s} = \gamma_* \left( \frac{\partial}{\partial s} \right), \quad (2a)$$

$$\gamma' := v^a \frac{\partial \gamma}{\partial \sigma^a} = \gamma_* \left( v^a \frac{\partial}{\partial \sigma^a} \right) = v^a \gamma_* \left( \frac{\partial}{\partial \sigma^a} \right). \quad (2b)$$

That the integral lines of  $s \mapsto \gamma(s, \vec{\sigma})$  be geodesics means that

$$\nabla_{\dot{\gamma}} \dot{\gamma} = 0. \quad (3)$$

Show that

$$L_{\dot{\gamma}} \gamma' = 0. \quad (4)$$

Next consider the projection of  $\gamma'$  orthogonal to  $\dot{\gamma}$ :

$$\gamma'_\perp := \gamma' - g(\gamma', \dot{\gamma}) \dot{\gamma}. \quad (5)$$

Show that

$$L_{\dot{\gamma}} \gamma'_\perp = 0. \quad (6)$$

(Hint: (6) follows from (5) once you prove  $\dot{\gamma}(g(\dot{\gamma}, \gamma')) = 0$ . To show that, use covariant derivatives, the geodesic equation (3), the fact that (5) implies  $\nabla_{\dot{\gamma}} \gamma' = \nabla_{\gamma'} \dot{\gamma}$  (for torsion-free connection), and that  $\nabla_{\gamma'} \dot{\gamma}$  is perpendicular to  $\dot{\gamma}$ .)

Finally prove the so-called *Jacobi-Equation of Geodesic Deviation*:

$$\nabla_{\dot{\gamma}} \nabla_{\dot{\gamma}} \gamma'_\perp = R(\dot{\gamma}, \gamma'_\perp) \dot{\gamma}. \quad (7)$$

(Hint: Recall the definition of curvature:  $(\nabla_X \nabla_Y - \nabla_Y \nabla_X - \nabla_{[X, Y]})Z = R(X, Y)Z$  and use it for  $X = Z = \dot{\gamma}$  and  $Y = \gamma'_\perp$ . Recall also that the connection is torsion-free.)

### Problem 2

The purpose of this exercise is to express equation (7) along the integral line of a single geodesic  $s \mapsto \gamma(s, \vec{\sigma})$  for some fixed value  $\vec{\sigma}$ , in terms of components of a conveniently chosen basis. Use an adapted orthonormal  $\{e_\alpha : \alpha = 0, 1, 2, 3\}$ , i.e.  $e_0 = \dot{\gamma}$  and  $g(e_\alpha, e_\beta) = \eta_{\alpha\beta} = \text{diag}(1, -1, -1, -1)$ , which is covariant constant along the selected integral line, i.e.  $\nabla_{\dot{\gamma}} e_\alpha = 0$  holds along that integral line. Show that (7) is then equivalent to

$$\frac{d^2 n^a}{ds^2} = R^a{}_{00b} n^b, \quad (8)$$

where we set  $n = n^a e_a := \gamma'_\perp$ . How do you interpret this equation?

### Problem 3

Write down the Jacobi equation (8) for the metric

$$g = \left(1 + \frac{2\phi(\vec{x})}{c^2}\right) c dt \otimes c dt - \left(1 - \frac{2\phi(\vec{x})}{c^2}\right) d\vec{x} \otimes d\vec{x} \quad (9)$$

in linear order in  $\phi/c^2$  (Newtonian approximation) and for a geodesic congruence for which at time  $t = 0$  all geodesics are running parallel to the  $t$  axis (all particles are rest at this moment of time). What is this equation telling you?

### Problem 4

Calculate all curvature components to linear order of a linearised, plane gravitational wave (compare Problem 3 of Sheet 7) in (+)-mode:

$$\begin{aligned} g = & c dt \otimes c dt - (1 - h_+(z - ct)) dx \otimes dx \\ & - (1 + h_+(z - ct)) dy \otimes dy \\ & - dz \otimes dz. \end{aligned} \quad (10)$$

Write down the Jacobi equation (8) for the timelike geodesic congruence of the metric (10) in which the spatial coordinates are all constant along each geodesic. (Hence the parameters  $\vec{\sigma}$  may be identified with the coordinates  $\vec{x} = (x, y, z)$ .) Again: What is this equation telling you?

### Problem 5 (for the Bravehearts)

Calculate all Riemann-, Ricci, and Einstein-Tensor components of the most general spherically symmetric metric ( $R$  is a constant of dimension “length”)

$$\begin{aligned} g = & e^{2a(t,r)} dx^0 \otimes dx^0 \\ & - e^{2b(t,r)} dr \otimes dr \\ & - e^{2c(t,r)} R^2 (d\theta \otimes d\theta + \sin^2(\theta) d\varphi \otimes d\varphi) \end{aligned} \quad (11)$$

with respect to the orthonormal co-frame  $\theta^0 = e^a dx^0$ ,  $\theta^1 = e^b dr$ ,  $\theta^2 = e^c R d\theta$ , and  $\theta^3 = e^c R \sin(\theta) d\varphi$ .