

Sheet 4: SolutionsProblem 1

$$\sigma^\alpha = (E_2, \vec{\sigma}) \quad (4.1.1)$$

$$\tilde{\sigma}^\alpha = (E_2, -\vec{\sigma}) \quad (4.1.2)$$

$$\sigma_\alpha := \eta_{\alpha\beta} \sigma^\beta = (E_2, -\vec{\sigma}) \quad (4.1.3)$$

$$\tilde{\sigma}_\alpha := \eta_{\alpha\beta} \tilde{\sigma}^\beta = (E_2, \vec{\sigma}) \quad (4.1.4)$$

1)

We know from Problem 1 of Sheet 2 that

$$M = 2 \text{Trace}(M) E_2 - \underbrace{\sigma_a M \sigma^a}_{\delta^{ab} \sigma_a M \sigma_b}$$

Hence

$$\underbrace{M + \delta^{ab} \sigma_a M \sigma_b}_{\sigma_\alpha M \tilde{\sigma}^\alpha} = 2 \text{Trace}(M) E_2 \quad (4.1.5)$$

2)

$$\sigma_\alpha \tilde{\sigma}^\beta = \frac{1}{2} (\sigma_\alpha \tilde{\sigma}^\beta + \sigma_\beta \tilde{\sigma}^\alpha) \quad (4.1.6)$$

$$a) \quad \alpha = 0, \beta = 0 \Rightarrow \frac{1}{2} (2 \sigma_0 \tilde{\sigma}_0) = E_2 \quad (4.1.7)$$

$$b) \quad \alpha = 0, \beta = b \Rightarrow \frac{1}{2} (\sigma_0 \tilde{\sigma}^b + \sigma_b \tilde{\sigma}_0) \\ = \frac{1}{2} (\sigma^b - \sigma^b) = 0 \quad (4.1.8)$$

$$\begin{aligned}
 \text{c) } \alpha = a, \beta = b &\leadsto \frac{1}{2} (\sigma_a \tilde{\sigma}_b + \sigma_b \tilde{\sigma}_a) \\
 &= -\frac{1}{2} (\sigma^a \sigma^b + \sigma^b \sigma^a) \\
 &= -\delta^{ab} E_2
 \end{aligned} \tag{4.1.9}$$

$$(\text{Recall: } \sigma^a \sigma^b = \delta^{ab} E_2 + i \varepsilon^{abc} \sigma^c)$$

Hence:

$$\sigma(\alpha \sigma \beta) = \eta_{\alpha\beta} E_2. \tag{4.1.10}$$

3) For any 4-tuple $X^\alpha \in \mathbb{R}^4$ have $X^\alpha \sigma_\alpha \in \text{Hermitean } (2 \times 2)\text{-matrices}$.
 If $X^\alpha = (X^0, X^1, X^2, X^3) = (X^0, \vec{X})$ and $\sigma_\alpha = (E_2, -\vec{\sigma})$ then

$$\begin{aligned}
 X^\alpha \sigma_\alpha &= X^0 \sigma_0 - \vec{X} \cdot \vec{\sigma} \\
 &= \begin{pmatrix} X^0 - X^3 & -X^1 + iX^2 \\ -X^1 - iX^2 & X^0 + X^3 \end{pmatrix}
 \end{aligned} \tag{4.1.11}$$

Note

$$\begin{aligned}
 \det(X^\alpha \sigma_\alpha) &= (X^0)^2 - \vec{X}^2 \\
 &= \eta_{\alpha\beta} X^\alpha X^\beta.
 \end{aligned} \tag{4.1.12}$$

The map $\sigma: \mathbb{R}^4 \rightarrow \text{Herm}(2 \times 2)$,
 $X^\alpha \mapsto X^\alpha \sigma_\alpha$ (4.1.13)

is a bijection (since clearly injective and linear map between real vector-spaces of the same dimension).

So if $h \in \text{Herm}(2 \times 2)$ then the inverse map to σ is

$$\begin{aligned} \sigma^{-1} : \text{Herm}(2 \times 2) &\rightarrow \mathbb{R}^4 \\ h &\mapsto X^\alpha = \frac{1}{2} \text{Trace}(h \tilde{\sigma}^\alpha). \end{aligned} \quad (4.1.14)$$

This is due to (4.1.10): if $h = X^\beta \sigma_\beta$, then

$$\begin{aligned} \text{Trace}(h \tilde{\sigma}^\alpha) &= \eta^{\alpha\gamma} \text{Trace}(h \tilde{\sigma}_\gamma) \\ &= \eta^{\alpha\gamma} \text{Trace}(X^\beta \sigma_\beta \tilde{\sigma}_\gamma) \\ &= \eta^{\alpha\gamma} X^\beta \text{Trace}(\sigma_\beta \tilde{\sigma}_\gamma) \\ &= \eta^{\alpha\gamma} X^\beta \text{Trace}(\sigma_{(\beta} \tilde{\sigma}_{\gamma)}) \\ &= \eta^{\alpha\gamma} \eta_{\beta\gamma} X^\beta \text{Trace}(E_2) \\ &\stackrel{(4.1.10)}{\rightarrow} \\ &= 2 X^\alpha. \end{aligned} \quad (4.1.15)$$

$$\begin{aligned} \leadsto X^\alpha &= \frac{1}{2} \text{Trace}(h \tilde{\sigma}^\alpha). \\ &= \frac{1}{2} \text{Trace}(\tilde{\sigma}^\alpha h) \end{aligned} \quad (4.1.16)$$

As $A(X^\alpha \sigma_\alpha) A^\dagger$ is again hermitian, it must be of the form $X'^\alpha \sigma_\alpha$, where X'^α depends linearly on X ; hence

$$X'^\alpha = L^\alpha_\beta X^\beta \quad (4.1.17)$$

$$\text{or } A(X^\alpha \sigma_\alpha) A^\dagger = L^\alpha_\beta X^\beta \sigma_\alpha \quad (4.1.18)$$

(1) As (4.1.18) holds identically for all $X^a \in \mathbb{R}^4$, we get

$$A \sigma_\beta A^+ = L^a{}_\beta \sigma_a \quad (4.1.19)$$

Using (4.1.16) we get

$$L^a{}_\beta = \frac{1}{2} \text{Trace}(\tilde{\sigma}^a A \sigma_\beta A^+) \quad (4.1.20)$$

4) $\det (A (X^a \sigma_a) A^+)$

$$= \det(A) \cdot \overline{\det(A)} \det(X^a \sigma_a)$$

$\begin{array}{ccc} \text{"} & \text{"} & \text{"} \\ 1 & 1 & \eta(X, X) \end{array}$

$$= \det((LX)^a \sigma_a) = \eta(LX, LX)$$

\hookrightarrow (4.1.18)

$$\Rightarrow \eta(LX, LX) = \eta(X, X) \quad (4.1.21)$$

$$\Leftrightarrow L \in O(\mathbb{R}^4, \eta). \quad (4.1.22)$$

Moreover, from (4.1.20) get

$$\begin{aligned} L^0{}_0 &= \frac{1}{2} \text{Trace} \left(\underset{\substack{\text{"} \\ \mathbb{E}_2}}{\tilde{\sigma}^0} A \underset{\substack{\text{"} \\ \mathbb{E}_2}}{\sigma_0} A^+ \right) \\ &= \frac{1}{2} \text{Trace}(A A^+) > 0 \end{aligned} \quad (4.1.23)$$

$\Rightarrow L$ is time-orientation preserving.

Similarly, again from (4.1.20) get

$$L^{\alpha} = \text{Trace}(L)$$

$$= \frac{1}{2} \text{Trace}(\tilde{\sigma}^{\alpha} A \sigma_{\alpha} A^{\dagger})$$

$$= \frac{1}{2} \text{Trace}(A \sigma_{\alpha} A^{\dagger} \tilde{\sigma}^{\alpha})$$

$$= \text{Trace}(A \text{Trace}(A^{\dagger}) E_2)$$

↳ using (4.1.5)

$$= \text{Trace}(A) \text{Trace}(A^{\dagger})$$

$$= |\text{Trace}(A)|^2 \geq 0$$

(4.1.24)

A space-reflection ("parity-transp.")
would be given by

$$\{L^{\alpha} \beta\} = \begin{pmatrix} 1 & \vec{0}^T \\ \vec{0} & -E_3 \end{pmatrix}$$

(4.1.25)

with $\text{Trace}(L) = -2$. Hence
 $L = p(A)$ cannot be a space-reflection.

$\Rightarrow p: SL(2, \mathbb{C}) \rightarrow O(\mathbb{R}^4, \eta)$ contains
neither time- nor space-reflections

$$\Rightarrow \text{Image}(p) \subseteq \text{Lor}_+^{\uparrow}$$

(4.1.26)

5) Multiplication of (4.1.19) with $\tilde{\sigma}^\beta$ gives, using (4.1.5):

$$\begin{aligned} \sigma_\alpha L^\alpha_\beta \tilde{\sigma}^\beta &= A \sigma_\beta A^\dagger \tilde{\sigma}^\beta \\ &= 2 \operatorname{Trace}(A^\dagger). \end{aligned} \quad (4.1.27)$$

Taking the determinant of both sides gives, due to $\det(A) = 1$,

$$\begin{aligned} 4 [\operatorname{Trace}(A^\dagger)]^2 &= \det(\sigma_\alpha L^\alpha_\beta \tilde{\sigma}^\beta) \\ \text{or } 2 \operatorname{Trace}(A^\dagger) &= \pm \sqrt{\det(\sigma_\alpha L^\alpha_\beta \tilde{\sigma}^\beta)} \end{aligned} \quad (4.1.28)$$

Re-inserting this into (4.1.27) gives

$$A = \pm \frac{\sigma_\alpha L^\alpha_\beta \tilde{\sigma}^\beta}{\sqrt{\det(\sigma_\alpha L^\alpha_\beta \tilde{\sigma}^\beta)}} \quad (4.1.28)$$

Finally we have that the map

$$\begin{aligned} p: SL(2, \mathbb{C}) &\rightarrow \operatorname{Lor}^{\uparrow}_+ \\ A &\mapsto p(A) = \{L^\alpha_\beta\} \end{aligned} \quad (4.1.29)$$

is a group homomorphism (as we should have noted already in part 4); this follows from the following: Let $A, B \in SL(2, \mathbb{C})$ with $p(A) = L$ and $p(B) = M$; that is

$$A \sigma_{\beta} A^{\dagger} = L^{\alpha} \beta \sigma_{\alpha} \quad (4.1.30)$$

$$B \sigma_{\beta} B^{\dagger} = M^{\gamma} \beta \sigma_{\gamma} \quad (4.1.31)$$

Then

$$\begin{aligned} & AB \sigma_{\beta} (AB)^{\dagger} \\ &= A B \sigma_{\beta} B^{\dagger} A^{\dagger} \\ &= A M^{\gamma} \beta \sigma_{\gamma} A^{\dagger} \\ &= M^{\gamma} \beta A \sigma_{\gamma} A^{\dagger} \\ &= M^{\gamma} \beta L^{\alpha} \sigma_{\alpha} \\ &= (L \cdot M)^{\alpha} \beta \sigma_{\alpha} \end{aligned} \quad (4.1.32)$$

$$\text{Hence } p(AB) = p(A) \cdot p(B) \quad (4.1.33)$$

$$\text{Also } p(E_2) = E_4. \quad (4.1.34)$$

$\Rightarrow p$ is group homomorphism.

From (4.1.28) it follows that for any $\{L^{\alpha} \beta\}$ in a neighbourhood of E_4 there is an A s.t. $p(A) = L$. Hence p maps onto an open neighbourhood of the identity in Lor^{\uparrow}_+ . Being a group homomorphism, it follows that p is surjective (since any element L of Lor^{\uparrow}_+ is the finite product of

elements from the open neighbourhood
(we have seen that this is true for any
connected topological group).

The kernel of $p: SL(2, \mathbb{C}) \rightarrow \text{Lo}_+^{\uparrow}$
is given by those A for which

$$A \sigma_{\alpha} A^{\dagger} = \sigma_{\alpha} \quad \forall \alpha \in \{0, 1, 2, 3\} \quad (4.1.35)$$

For $\alpha = 0$ this implies $A A^{\dagger} = E_2$
 $\Leftrightarrow A^{\dagger} = A^{-1}$; hence

$$A \sigma_{\alpha} = \sigma_{\alpha} A \quad (4.1.36)$$

$\Leftrightarrow A$ commutes with all complex
 2×2 -matrices (since $\text{Span}_{\mathbb{C}} \{\sigma^{\alpha}\} =$
 $\text{End}(\mathbb{C}^2)$) $\Leftrightarrow A = a E_2$, $a \in \mathbb{C}$,
with $a = \pm 1$ since $\det(A) = 1$.

Hence

$$\text{kernel}(p) = \{\pm E_2\}. \quad (4.1.37)$$

Problem 2

$$[P_a, P_b] = 0 \quad (4.2.1)$$

$$[M_{ab}, P_c] = \eta_{bc} P_a - \eta_{ac} P_b \quad (4.2.2)$$

$$[M_{ab}, M_{cd}] = \eta_{ad} M_{bc} + \eta_{bc} M_{ad} \\ - \eta_{ac} M_{bd} - \eta_{bd} M_{ac} \quad (4.2.3)$$

1)

Contraction of (4.2.2) with η^{bc} gives

$$\eta^{bc} [M_{ab}, P_c] = n P_a - P_a \\ = (n-1) P_a. \quad (4.2.4)$$

Here $n = \text{dimension } V$. Hence

$$P_a = \frac{1}{n-1} \underbrace{\eta^{bc} [M_{ab}, P_c]}_{\text{Sum of commutators}} \quad (4.2.5)$$

Contraction of (4.2.3) with η^{ad} gives

$$\eta^{ad} [M_{ab}, M_{cd}] = n M_{bc} + \eta_{bc} \eta^{ad} M_{ad} \\ - M_{bc} - M_{bc} \quad (4.2.6)$$

where $\eta^{ad} M_{ad} = 0$ because $\eta^{ad} = \eta^{da}$ and $M_{ad} = -M_{da}$. Hence

$$M_{bc} = \frac{1}{n-2} \eta^{ad} [M_{ab}, M_{cd}] \quad (4.2.7)$$

where the right-hand side is again a sum of Lie-brackets.

As every basis vector P_a, M_{bc} is the sum of Lie-brackets, so is any element of the Lie-algebra, which is hence perfect.

2.)

Let $\rho : \text{Lie}(V \times O(V, \eta)) \rightarrow \text{End}(W)$ a representation in the one-dim. vector space W ; $\dim_{\mathbb{F}}(W) = 1$. Then all $\rho(X) \in \text{End}(W)$ commute, i.e.

$$\rho(X)\rho(Y) - \rho(Y)\rho(X) = 0 \quad (4.2.8)$$

The left hand side equals $\rho([X, Y])$. Hence the ρ -image of a commutator is zero. Since the Lie-algebra is perfect

$$\rho \equiv 0 \quad (4.2.9)$$

This proves that a one-dimensional representation of a perfect Lie-algebra is trivial (i.e. identically zero).

In contrast, the abelian subgroup of translations has many non-trivial one-dimensional representations. For example, let $k \in V^*$, then

$$\mathfrak{g}(v) = \mathfrak{g}_k(v) := k(v) \in \mathbb{F} \quad (4.2.10)$$

is such a representation. If $\mathbb{F} = \mathbb{C}$
we may choose

$$k = i \eta(p, \cdot) \quad (4.2.11)$$

so that

$$\mathfrak{g}(v) = i \eta(p, v) = i p \cdot v \in i \mathbb{R} \quad (4.2.12)$$

corresponding to the unitary
representation

$$D(v) = \exp(i p \cdot v) \in \mathbb{C}, \quad (4.2.13)$$

i.e.

$$\dot{D} = \mathfrak{g}_k \quad (4.2.14)$$

It is remarkable that none of these
1-dim. rep. of V extend to
 $V \times \mathcal{O}(V, \eta)$.

$$3.) \quad P_a \mapsto \bar{P}_a := P_a \quad (4.2.15)$$

$$M_{ab} \mapsto M'_{ab} := M_{ab} - (X_a P_b - X_b P_a) \quad (4.2.16)$$

where $(X^1, \dots, X^n) \in \mathbb{F}^n$. It immediately
follows that

$$[\bar{P}_a, \bar{P}_b] = 0 \quad (4.2.17)$$

$$[\bar{M}_{ab}, \bar{P}_c] =$$

$$[M_{ab} - (X_a P_b - X_b P_a), P_c]$$

$$= [M_{ab}, P_c]$$

$$= \eta_{bc} P_a - \eta_{ac} P_b$$

$$= \eta_{bc} \bar{P}_a - \eta_{ac} \bar{P}_b$$

(4.2.18)

$$[\bar{M}_{ab}, \bar{M}_{cd}] =$$

$$[M_{ab} - (X_a P_b - X_b P_a), M_{cd} - (X_c P_d - X_d P_c)]$$

$$= [M_{ab}, M_{cd}] - [X_a P_b - X_b P_a, M_{cd}]$$

$$- [M_{ab}, X_c P_d - X_d P_c]$$

$$= \eta_{ad} M_{bc} + \eta_{bc} M_{ad} - \eta_{ac} M_{bd} - \eta_{bd} M_{ac}$$

$$- X_a [P_b, M_{cd}] + X_b [P_a, M_{cd}]$$

$$- X_c [M_{ab}, P_d] + X_d [M_{ab}, P_c]$$

$$= \eta_{ad} \bar{M}_{bc} + \eta_{bc} \bar{M}_{ad} - \eta_{ac} \bar{M}_{bd} - \eta_{bd} \bar{M}_{ac} \quad (4.2.19)$$

$$+ \underbrace{\eta_{ad} (X_b P_c - X_c P_b)}_1 + \underbrace{\eta_{bc} (X_a P_d - X_d P_a)}_2$$

$$- \underbrace{\eta_{ac} (X_b P_d - X_d P_b)}_3 - \underbrace{\eta_{bd} (X_a P_c - X_c P_a)}_4$$

$$+ X_a (\underbrace{\eta_{bd} P_c}_4 - \underbrace{\eta_{bc} P_d}_2) - X_b (\underbrace{\eta_{ad} P_c}_1 - \underbrace{\eta_{ac} P_d}_3)$$

$$- X_c (\underbrace{\eta_{bd} P_a}_4 - \underbrace{\eta_{ad} P_b}_1) + X_d (\underbrace{\eta_{bc} P_a}_2 - \underbrace{\eta_{ac} P_b}_3)$$

These terms cancel

Problem 3

$$T_{(a, A)} \varphi = \varphi \circ \phi^{-1}_{(a, A)} \quad (4.3.1)$$

$$\begin{aligned} \text{where } \phi : ((a, A), V) &\mapsto \phi_{(a, A)}(V) \\ &=: AV + a \end{aligned} \quad (4.3.2)$$

is an action

$$\begin{aligned} 1.) \quad T_{(a, A)} \circ T_{(b, B)} \varphi &= (T_{(b, B)} \varphi) \circ \phi^{-1}_{(a, A)} \\ &= \varphi \circ \phi^{-1}_{(b, B)} \circ \phi^{-1}_{(a, A)} \\ &= \varphi \circ (\phi_{(a, A)} \circ \phi_{(b, B)})^{-1} \\ &= \varphi \circ \phi^{-1}_{(a + Ab, AB)} \\ &= T_{(a + Ab, AB)} \varphi. \end{aligned} \quad (4.3.3)$$

$$\text{And clearly also } T \text{id}_V = \text{id} |_{C^\infty(V, \mathbb{R})}. \quad (4.3.4)$$

2) Let $S \mapsto (a(s), A(s))$ be curve in G with $a(0) = 0$, $A(0) = \text{id}_V$

$$\begin{aligned} &(T_{(a(s), A(s))} \varphi)(X) \\ &= \varphi(\phi^{-1}_{(a(s), A(s))} X) = \varphi(\phi_{(a(s), A(s))}^{-1} X) \end{aligned} \quad (4.3.5)$$

With $(a, A)^{-1} = (-A^{-1}a, A^{-1})$ get

$$(T_{(a(s), A(s))} f)(x) = f(A^{-1}(s)(x - a(s)))$$

$$\frac{d}{ds} \Big|_{s=0} (A^{-1}(s)(x - a(s)))$$

$$= -\dot{A}x - \dot{a}$$

$$\Rightarrow \frac{d}{ds} \Big|_{s=0} (T_{(a(s), A(s))} f)(x)$$

$$= \underbrace{Df(x)}_{\text{derivative of } f \text{ at } x} (-\dot{a} - \dot{A}x)$$

(4.3.6)

derivative of f at x .

3) LM

$$P_a = (e_a, 0)$$

(4.3.7)

$$M_{ab} = (0, e_a \otimes \otimes b - e_b \otimes \otimes a)$$

(4.3.8)

basis of $\text{Lie}(V \times \mathcal{O}(V, \eta))$

With

$$\dot{T}_{(\dot{a}, \dot{A})} f(x) = Df(x) (-\dot{a} - \dot{A}x) \quad (4.3.9)$$

have

$$\dot{T}_{P_a} f(x) = Df(x) (-e_a)$$

$$= -De_a f(x) = -\partial_a f(x) \quad (4.3.10)$$

$$\begin{aligned}
 (\dot{T}_{M_{ab}} f)(x) &= \mathcal{D}f(x) (-M_{ab}x) \\
 &= \mathcal{D}f(x) (-e_a x_b + e_b x_a) \\
 &= x_a \partial_b f(x) - x_b \partial_a f(x). \quad (4.3.11)
 \end{aligned}$$

$$[\dot{T}_{P_a}, \dot{T}_{P_b}] = [-\partial_a, -\partial_b] = 0 \quad (4.3.12)$$

$$\begin{aligned}
 [\dot{T}_{M_{ab}}, \dot{T}_{P_c}] &= [x_a \partial_b - x_b \partial_a, -\partial_c] \\
 &= (\partial_c x_a) \partial_b - (\partial_c x_b) \partial_a \\
 &= \eta_{bc} \dot{T}_{P_a} - \eta_{ac} \dot{T}_{P_b} \quad (4.3.13)
 \end{aligned}$$

$$\begin{aligned}
 [\dot{T}_{M_{ab}}, \dot{T}_{M_{cd}}] &= \\
 &= [x_a \partial_b - x_b \partial_a, x_c \partial_d - x_d \partial_c] \\
 &= \underbrace{\eta_{bc} x_a \partial_d}_2 - \underbrace{\eta_{bd} x_a \partial_c}_4 \\
 &\quad - \underbrace{\eta_{ac} x_b \partial_d}_3 + \underbrace{\eta_{ad} x_b \partial_c}_1 \\
 &\quad - \underbrace{\eta_{ad} x_c \partial_b}_1 + \underbrace{\eta_{bd} x_c \partial_a}_4 \\
 &\quad + \underbrace{\eta_{ac} x_d \partial_b}_3 - \underbrace{\eta_{bc} x_d \partial_a}_2 \\
 &= \eta_{ad} \dot{T}_{M_{bc}} + \eta_{bc} \dot{T}_{M_{ad}} \\
 &\quad - \eta_{ac} \dot{T}_{M_{bd}} - \eta_{bd} \dot{T}_{M_{ac}} \quad (4.3.14)
 \end{aligned}$$

Problem 4

S4.16

We consider $\text{Lie}(O(V, \eta))$ with $V = 5$ -dim. real vector space with non-deg. symmetric bilinear form of signature $(1, -1, -1, -1, \sigma)$, where $\sigma = -1$ (de Sitter case) or $\sigma = 1$ (anti-de Sitter case). We know that a basis of $\text{Lie}(O(V, \eta))$ is given by the $10 = \frac{1}{2}5(5-1)$ element set

$$\{ M_{ab} = e_a \otimes \theta_b - e_b \otimes \theta_a \mid 0 \leq a < b \leq 5 \}, \quad (4.4.1)$$

with $\{ e_a \mid 0 \leq a \leq 4 \}$ basis of V , $\{ \theta^a \mid 0 \leq a \leq 4 \}$ basis of V^* dual to $\{ e_a \}$, i.e. $\theta^a(e_b) = \delta^a_b$, $\eta_{ab} := (e_a, e_b)$, and

$$\theta_a = \eta_{ab} \theta^b. \quad (4.4.2)$$

Then

$$\begin{aligned} [M_{ab}, M_{cd}] &= \eta_{ad} M_{bc} + \eta_{bc} M_{ad} \\ &\quad - \eta_{ac} M_{bd} - \eta_{bd} M_{ac} \end{aligned} \quad (4.4.3)$$

We write the 10 elements $M_{ab} \in V \otimes V^*$ in the following form, where now $1 \leq a, b, c \leq 3$

	#	54, 17
$D_a := \frac{1}{2} \epsilon_a{}^{bc} M_{bc}$	3	(4.4.4)
$K_a := M_{0a}$	3	(4.4.5)
$T_a := M_{a4}$	3	
$T_0 := M_{04}$	1	
	10	

The Lie-products between them can be conveniently organised in lexicographical order

	#	
$[D_a, D_b], [D_a, K_b], [D_a, T_b], [D_a, T_0]$		
(3) (9) (9) (3)		(24)
$[K_a, K_b], [K_a, T_b], [K_a, T_0]$		
(3) (9) (3)		(15)
$[T_a, T_b], [T_a, T_0]$		
(3) (3)		6

where the numbers in round brackets denote the numbers of independent relations. In total there are 45, which is obvious from the number 10 indep. Mats and $\frac{1}{2} 10(10-1) = 5 \cdot 9 = 45$ independent $[M_{ab}, M_{cd}]$.

The relations may be calculated using (4.4.3). To understand the structure, let us restrict to representative examples:

$$\begin{aligned}
 [D_1, D_2] &= [M_{23}, M_{31}] \\
 &= \cancel{\eta_{21}} M_{33} + \eta_{33} M_{21} - \cancel{\eta_{23}} M_{31} - \cancel{\eta_{31}} M_{23} \\
 &= -M_{21} = M_{12} = D_3 \quad (4.4.6)
 \end{aligned}$$

$$\begin{aligned}
 [D_1, K_2] &= [M_{23}, M_{02}] \\
 &= \eta_{22} M_{30} + \cancel{\eta_{30}} M_{22} - \cancel{\eta_{20}} M_{32} - \cancel{\eta_{32}} M_{20} \\
 &= -M_{30} = M_{03} = K_3 \quad (4.4.7)
 \end{aligned}$$

$$\begin{aligned}
 [D_1, T_2] &= [M_{23}, M_{24}] \\
 &= \cancel{\eta_{24}} M_{32} + \cancel{\eta_{32}} M_{24} - \eta_{22} M_{34} - \cancel{\eta_{34}} M_{22} \\
 &= M_{34} = T_3 \quad (4.4.8)
 \end{aligned}$$

$$\begin{aligned}
 [D_1, T_0] &= [M_{23}, M_{04}] \\
 &= \cancel{\eta_{24}} M_{30} + \cancel{\eta_{30}} M_{24} - \cancel{\eta_{20}} M_{34} - \cancel{\eta_{34}} M_{20} \\
 &= 0 \quad (4.4.9)
 \end{aligned}$$

$$\begin{aligned}
 [K_1, K_2] &= [M_{01}, M_{02}] \\
 &= \eta_{02} M_{10} + \cancel{\eta_{10}} M_{02} - \eta_{00} M_{12} - \cancel{\eta_{12}} M_{00} \\
 &= -M_{12} = -D_3 \quad (4.4.10)
 \end{aligned}$$

$$[K_1, T_2] = [M_{01}, M_{24}]$$

$$= \cancel{\gamma_{04}} M_{12} + \cancel{\gamma_{12}} M_{04} - \cancel{\gamma_{02}} M_{14} - \cancel{\gamma_{14}} M_{02}$$

$$= 0$$

(4.4.11)

$$[K_1, T_1] = [M_{01}, M_{14}]$$

$$= \cancel{\gamma_{04}} M_{11} + \cancel{\gamma_{11}} M_{04} - \cancel{\gamma_{01}} M_{14} - \cancel{\gamma_{14}} M_{01}$$

$$= -M_{04} = -T_0$$

(4.4.12)

$$[K_1, T_0] = [M_{01}, M_{04}]$$

$$= \cancel{\gamma_{04}} M_{10} + \cancel{\gamma_{10}} M_{04} - \cancel{\gamma_{00}} M_{14} - \cancel{\gamma_{14}} M_{00}$$

$$= -M_{14} = -T_1$$

(4.4.13)

$$[T_1, T_2] = [M_{14}, M_{24}]$$

$$= \cancel{\gamma_{14}} M_{42} + \cancel{\gamma_{42}} M_{14} - \cancel{\gamma_{12}} M_{44} - \cancel{\gamma_{44}} M_{12}$$

$$= -\varepsilon M_{12} = -\sigma D_3$$

(4.4.14)

$$[T_1, T_0] = [M_{14}, M_{04}]$$

$$= \cancel{\gamma_{14}} M_{40} + \cancel{\gamma_{40}} M_{14} - \cancel{\gamma_{10}} M_{44} - \cancel{\gamma_{44}} M_{10}$$

$$= -\sigma M_{10} = \sigma M_{01} = \sigma K_1$$

(4.4.15)

In total we get

$$[D_a, D_b] = \epsilon_{ab}{}^c D_c \quad (4.4.16)$$

$$[D_a, K_b] = \epsilon_{ab}{}^c K_c \quad (4.4.17)$$

$$[D_a, T_b] = \epsilon_{ab}{}^c T_c \quad (4.4.18)$$

$$[D_a, T_0] = 0 \quad (4.4.19)$$

$$[K_a, K_b] = -\epsilon_{ab}{}^c D_c \quad (4.4.20)$$

$$[K_a, T_b] = -\delta_{ab} T_0 \quad (4.4.21)$$

$$[K_a, T_0] = -T_a \quad (4.4.22)$$

$$[T_a, T_b] = -\sigma \epsilon_{ab}{}^c D_c \quad (4.4.23)$$

$$[T_a, T_0] = \sigma K_a \quad (4.4.24)$$

2) Consider the linear maps defined by

$$\Pi : (T_0, T_a, K_a, D_a) \mapsto (T_0, -T_a, -K_a, D_a) \quad (4.4.25)$$

$$\Theta : (T_0, T_a, K_a, D_a) \mapsto (-T_0, T_a, -K_a, D_a) \quad (4.4.26)$$

$$\Gamma : (T_0, T_a, K_a, D_a) \mapsto (-T_0, -T_a, K_a, D_a) \quad (4.4.27)$$

Have

$$\Gamma = \Theta \circ \Pi = \Pi \circ \Theta \quad (4.4.28)$$

Π = space-orientation reversal

Θ = time-orientation reversal

That π and θ (and hence their composition Γ) are Lie-automorphisms means that, e.g.,

$$\pi[X, Y] = [\pi(X), \pi(Y)]$$

It follows by direct inspection that changing the signs of T_a and K_b or T_a and K_a simultaneously leaves (4.4.16-24) invariant, so that π and θ are indeed automorphisms.

$$3) \quad U_1 := \text{Span}\{D_a, K_b \mid 1 \leq a, b \leq 3\}$$

$$U_2 := \text{Span}\{D_a, T_b \mid 1 \leq a, b \leq 3\}$$

are obviously Lie-subalgebras isomorphic to $\text{Lie}(SL(2, \mathbb{C})) = \text{Lie}(Lor)$.

a) Contraction over U_1 means to set all commutators among T_a, T_b to zero and also all terms $\sim D_a, K_a$ between elements of U_1 and $\{T_a, T_b\}$. We then get. This amounts to setting the right-hand sides of (4.4.23-24) to zero, thereby yielding the Lie-algebra of the Poincaré group

Independent of σ .

(In comparison with relations of Skript, p. 3.31, set $D_a = \dot{R}_a$, $K_a = \dot{B}_a$, $T_a = \dot{P}_a$, $T_0 = -\dot{P}_0$).

b.) contraction over U_2 means to set all commutators among elements $\{K_a, T_a\}$ to zero and cancel all terms $\sim D_a, T_b$ in relations between U_2 and $\{K_a, T_a\}$

$$[D_a, D_b] = \epsilon_{ab}^c D_c$$

$$[D_a, K_b] = \epsilon_{ab}^c K_c$$

$$[D_a, T_b] = \epsilon_{ab}^c T_c$$

$$[D_a, T_0] = 0$$

$$[K_a, K_b] = 0$$

$$[K_a, T_b] = -\delta_{ab} T_0$$

$$[K_a, T_0] = 0$$

$$[T_a, T_b] = -\sigma \epsilon_{ab}^c D_c$$

$$[T_a, T_0] = \sigma K_a$$

That is, in contrast to the previous case, we set the right-hand sides to (4.4.20, 22) to zero. This implies that the ensuing Lie-algebra for $\sigma = 1$ is, in fact isomorphic to that of case a), as one sees by replacing $T_a \rightarrow K_a$, $K_a \rightarrow -T_a$:

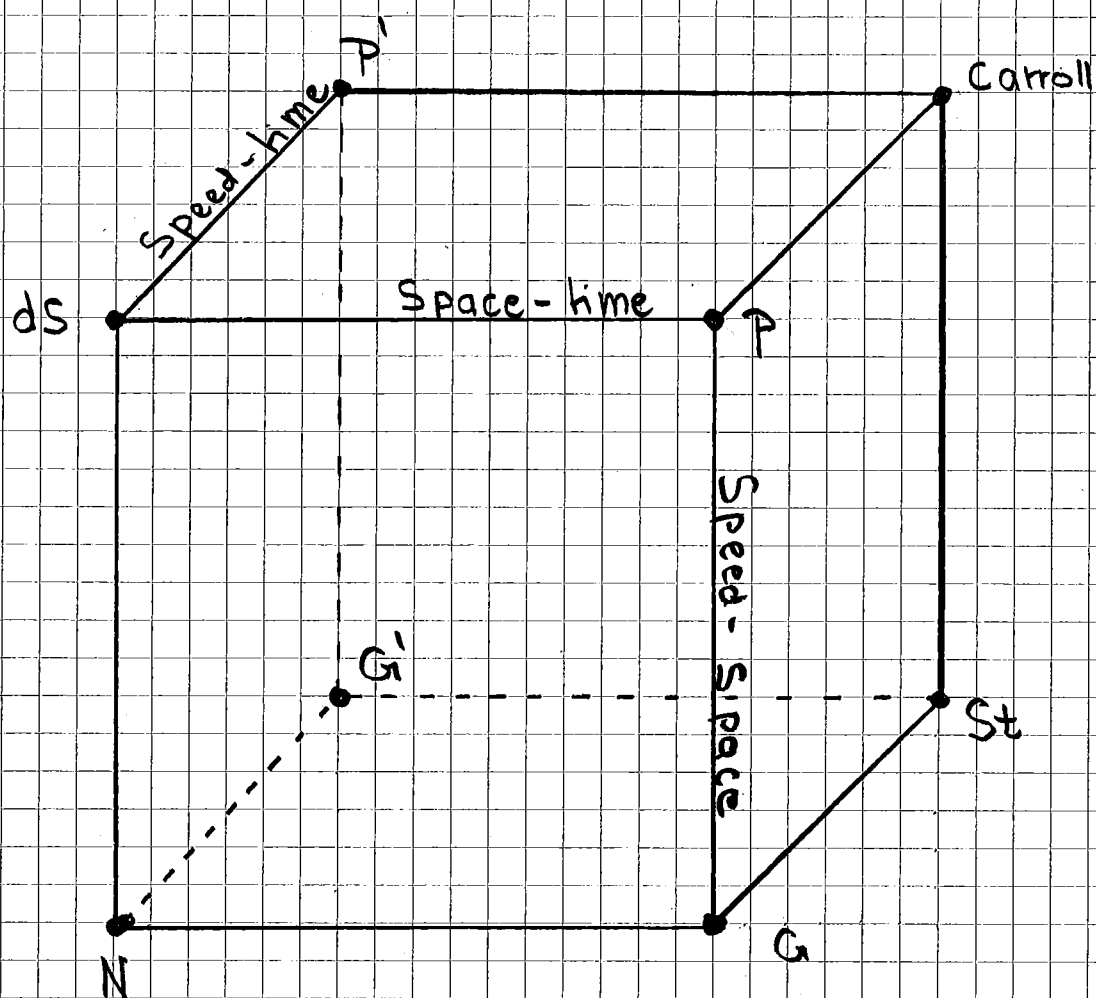
	over $\{D_a, K_b\}$ Small $\{T_a, T_b\}$	over $\{D_a, T_b\}$ Small $\{K_a, T_a\}$	1 & 2 2 & 1	1 & 8 Small $\{K_a, T_a\}$	2 & 8 Small $\{K_a, T_a\}$
$[D_a, D_b] = \epsilon_{ab}{}^c D_c$	"	"	"	"	"
$[D_a, K_b] = \epsilon_{ab}{}^c K_c$	"	"	"	"	"
$[D_a, T_b] = \epsilon_{ab}{}^c T_c$	"	"	"	"	"
$[D_a, T_0] = 0$	"	"	"	"	"
$[K_a, K_b] = -\epsilon_{ab}{}^c D_c$	"	0	0	0	0
$[K_a, T_b] = -\delta_{ab} T_0$	"	"	"	0	0
$[K_a, T_0] = -T_a$	"	0	0	"	0
$[T_a, T_b] = -\sigma \epsilon_{ab}{}^c D_c$	0	0	0	0	0
$[T_a, T_0] = \sigma K_a$	0	0	0	0	"

Poincaré para Poincaré

Galilei

para Galilei

$\sigma = +1, T_a \rightarrow K_a, K_a \rightarrow -T_a$



The contractions indicated here are

space-time : $\{T_a, T_b\} \rightarrow$ small

Speed-space : $\{K_a, T_b\} \rightarrow$ small

Speed-time : $\{K_a, T_b\} \rightarrow$ small

These lead to $P =$ "Poincaré",

$P' =$ "para Poincaré", $G =$ "Galilei",

$G' =$ "para Galilei", $C =$ "Carroll",

$N =$ "non-relativistic cosmological", and

$St =$ "static" Lie-algebras

For more discussion of these 7 contractions from $dS = de$ Sitter algebra, see the following paper:

Henri Bacry & Jean-Marc Lévy-Leblond:
"Possible Kinematics"

Journal of Mathematical Physics,
Volume 9, Number 10, year 1968,
pages 1605 - 1614.