# Exercises for the lecture on

## Foundations and Applications of Special Relativity

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### Sheet 3

# **Problem 1**

In the lecture we wrote down a Lorentz boosts with respect to inertial coordinates (ct, x, y, z). It reads (interpreted actively)

$$\begin{pmatrix} ct \\ x \\ y \\ z \end{pmatrix} \longmapsto B(\beta \mathbf{e}_x) \begin{pmatrix} ct \\ x \\ y \\ z \end{pmatrix}$$
(1a)

where

$$B(\beta \mathbf{e}_{x}) := \begin{pmatrix} \gamma & \beta \gamma & 0 & 0\\ \beta \gamma & \gamma & 0 & 0\\ 0 & 0 & 1 & 0\\ 0 & 0 & 0 & 1 \end{pmatrix}$$
(1b)

and

$$\gamma := \frac{1}{\sqrt{1 - \beta^2}}, \qquad \beta := \frac{\nu}{c}. \tag{1c}$$

1. Show that  $B(\beta e_x)$  is a symmetric and positive-definite linear map in  $\mathbb{R}^4$  with respect to the standard euclidean inner product, the eigenvalues of which are

$$\lambda_{1,2} = \gamma \pm \sqrt{\gamma^2 - 1} = \sqrt{\frac{1 \pm \beta}{1 \mp \beta}}, \qquad \lambda_{3,4} = 1, \qquad (2)$$

(note that  $\lambda_2 = \lambda_1^{-1}$ ) and the normalised eigenvectors of which are

$$\mathbf{E}_{1,2} = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 \\ \pm 1 \\ 0 \\ 0 \end{pmatrix}, \qquad \mathbf{E}_{3} = \begin{pmatrix} 0 \\ 0 \\ 1 \\ 0 \end{pmatrix}, \qquad \mathbf{E}_{3} = \begin{pmatrix} 0 \\ 0 \\ 0 \\ 1 \end{pmatrix}. \tag{3}$$

2. A spatial rotation with  $D \in SO(3)$  is represented, again with respect to the standard basis in  $\mathbb{R}^4$ , by the  $(4 \times 4)$  matrix (written in 1+3 form)

$$\mathbf{R}(\mathbf{D}) = \begin{pmatrix} \mathbf{1} & \mathbf{0}^{\mathsf{T}} \\ \mathbf{0} & \mathbf{D} \end{pmatrix} \,. \tag{4}$$

Show that

$$R(D)B(\beta \mathbf{e}_{x})R(D^{-1}) = B(\beta D\mathbf{e}_{x}),$$
(5)

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and that, consequently, a boost in a general direction  $\mathbf{n} = D\mathbf{e}_{\chi}$  is given by

$$B(\beta \mathbf{n}) = \begin{pmatrix} \gamma & \beta \gamma \mathbf{n}^{\top} \\ \beta \gamma \mathbf{n} & E_3 + (\gamma - 1) \mathbf{n} \otimes \mathbf{n}^{\top} \end{pmatrix}.$$
(6)

Here  $E_3$  denotes the unit  $(3 \times 3)$  matrix and  $\mathbf{n} \otimes \mathbf{n}^{\top}$  is the  $(3 \times 3)$ -matrix with components  $n_a n_b$ . Hence a Lorentz boost in the direction  $\mathbf{n}$  can also be written

$$\operatorname{ct} \mapsto \gamma(\operatorname{ct} + \beta(\mathbf{n} \cdot \mathbf{x})),$$
 (7a)

$$\mathbf{x} \mapsto \mathbf{x} + (\gamma - 1)\mathbf{n}(\mathbf{n} \cdot \mathbf{x}) + \beta \gamma \operatorname{nct}$$
 (7b)

- 3. Show that (7) also follows directly from (1) by writing  $\mathbf{x} = \mathbf{x}_{\parallel} + \mathbf{x}_{\perp}$ , where  $\mathbf{x}_{\parallel}$  and  $\mathbf{x}_{\perp}$  correspond to the components parallel and perpendicular to the boost direction  $\mathbf{e}_{x}$ .
- 4. Show that a general Boost transformation leaves a 2-dimensional plane in  $\mathbb{R}^4$  pointwise fixed and transforms the orthogonal plane (with respect to the euclidean inner product) non-trivially into itself with only fixed-point being the origin. Give a geometric description of the orbits of that action.

## Problem 2

Let  $M \in GL(\mathbb{R}^n)$  be a general invertible  $(n \times n)$ -matrix with real entries. Show that it has a unique so-called "polar decomposition"

$$M = PO \tag{8}$$

into an orthogonal matrix  $O \in O(3)$  and a symmetric and positive-definite matrix P (written in the order just stated).

Hint: Consider  $MM^{\top}$  and prove that it is symmetric and positive definite. Argue that there is a positive square-root  $P := \sqrt{MM^{\top}}$  which is also symmetric and positive definite, so that  $P^2 = MM^{\top}$ . Then define  $O := P^{-1}M$  and prove that it is orthogonal. This proves existence. For uniqueness assume  $M = P_1O_1 = P_2O_2$  with  $P_{1,2}$  symmetric and positive-definite and  $O_{1,2}$  orthogonal. Show that then  $P_1 = P_2$  and  $O_1 = O_2$  follow.

#### Problem 3

This problem is partly a repetition of Problem 3 on Sheet 2 which we did not yet come to discuss so far. Let  $\mathbf{V} : \mathbb{R}^3 \to \mathbb{R}^3$  be a vector field.

1. Show that the equations

$$\boldsymbol{\nabla} \cdot \mathbf{V} = \mathbf{0} \quad \text{and} \quad \boldsymbol{\nabla} \times \mathbf{V} = \mathbf{0} \tag{9}$$

permit rotations as symmetries in the sense given on Sheet 2.

2. Derive the form of the most general rotation-invariant vector field and find amongst them all solutions to the first equation (9) and also all solutions to the second equation.