

Lecturer: Olaf Lechtenfeld

Assistant: Nicolas Eicke

Perturbation theory for interacting ϕ^4 scalar field theory

Consider a scalar field theory with quartic self-interaction, described by

$$\mathcal{L} = \frac{1}{2} \partial_\mu \phi \partial^\mu \phi - \frac{1}{2} m^2 \phi^2 - \frac{\lambda}{4!} \phi^4. \tag{1}$$

In quantum field theory, all n -point correlation functions can be encoded in a single object called the generating functional, $Z[J]$, as

$$G_n(x_1, x_2, \dots, x_n) = (-i)^n \left[\frac{\delta^n Z[J]}{\delta J(x_1) \delta J(x_2) \dots \delta J(x_n)} \right]_{J=0}. \tag{2}$$

For the ϕ^4 theory, the generating functional is given by

$$Z[J] = \frac{\exp \left(\frac{-i\lambda}{4!} \int d^4z \left(-i \frac{\delta}{\delta J(z)} \right)^4 \right) Z_0[J]}{\left\{ \exp \left(\frac{-i\lambda}{4!} \int d^4z \left(-i \frac{\delta}{\delta J(z)} \right)^4 \right) Z_0[J] \right\}_{J=0}}, \tag{3}$$

where Z_0 is the free generating functional

$$Z_0[J] = Z_0[0] \exp \left(-\frac{1}{2} \int d^4x d^4y J(x) D_F(x-y) J(y) \right). \tag{4}$$

Assuming small interaction coupling, $\lambda \ll 1$, we can use perturbation theory. The Feynman rules for the ϕ^4 theory read:

$$\begin{aligned} \bullet \text{---} \bullet &= D_F(x-y) \\ \bullet \text{---} \times &= \int d^4z D_F(x-z) iJ(z) \\ \times \text{---} \times &= \int d^4x \int d^4y iJ(x) D_F(x-y) iJ(y) \\ \diagdown \diagup &= \left(\frac{-i\lambda}{4!} \right) \int d^4z \end{aligned}$$

- 1 (a) Apply the functional derivative $-i\frac{\delta}{\delta J(z)}$ to $Z_0[J]$, and draw the corresponding diagram. What does $-i\frac{\delta}{\delta J(z)}$ do to a diagram?
- (b) By expanding $Z[J]$ to first order in λ , and applying $-i\frac{\delta}{\delta J(z)}$ four times, show that to $\mathcal{O}(\lambda)$

$$Z[J] = \frac{Z_0[J]}{\left[1 - \frac{i\lambda}{4!} \int d^4x (3(D_F(0))^2) Z_0[0]\right]} \times \left[1 - \frac{i\lambda}{4!} \int d^4x \left(3(D_F(0))^2 + 6D_F(0) \left(\int d^4y D_F(x-y)J(y)\right)^2 + \left(\int d^4y D_F(x-y)J(y)\right)^4\right)\right]. \quad (5)$$

You may choose whether to work explicitly or to use the Feynman diagrams.

- (c) Show that the vacuum diagrams, which diverge, cancel thanks to the normalization.

2 In terms of Feynman diagrams, the four-point function reads

$$G_4(x_1, x_2, x_3, x_4) = \left(\begin{array}{c} 1 \text{ --- } 2 \\ 3 \text{ --- } 4 \end{array} + \begin{array}{c} 1 \\ | \\ 3 \end{array} \begin{array}{c} 2 \\ | \\ 4 \end{array} + \begin{array}{c} 1 \quad 2 \\ \diagdown \quad \diagup \\ 3 \quad 4 \end{array} \right) \\ - \frac{i\lambda}{2} \left(\begin{array}{c} 1 \text{ --- } 2 \\ \text{loop} \\ 3 \text{ --- } 4 \end{array} + \begin{array}{c} 1 \text{ --- } 2 \\ \text{loop} \\ 3 \text{ --- } 4 \end{array} + \begin{array}{c} 1 \\ | \\ \text{loop} \\ | \\ 3 \end{array} \begin{array}{c} 2 \\ | \\ 4 \end{array} + \begin{array}{c} 1 \\ | \\ \text{loop} \\ | \\ 3 \end{array} \begin{array}{c} 2 \\ | \\ 4 \end{array} + \begin{array}{c} 1 \quad 2 \\ \diagdown \quad \diagup \\ \text{loop} \\ \diagdown \quad \diagup \\ 3 \quad 4 \end{array} + \begin{array}{c} 1 \quad 2 \\ \diagdown \quad \diagup \\ \text{loop} \\ \diagdown \quad \diagup \\ 3 \quad 4 \end{array} \right) \\ - i\lambda \begin{array}{c} 1 \quad 2 \\ \diagdown \quad \diagup \\ 3 \quad 4 \end{array} + \mathcal{O}(\lambda^2). \quad (6)$$

What diagrams appear in $G_4(x_1, x_2, x_3, x_4)$ at $\mathcal{O}(\lambda^2)$? Take the symmetry factors into account!

- 3 Only the connected Feynman diagrams in a correlation function contribute to the non-trivial (off-diagonal) part of the S -matrix. Show to $\mathcal{O}(\lambda)$ that the functional $W[J] = -i \ln Z[J]$ generates only the connected diagrams of $G_4(x_1, x_2, x_3, x_4)$.