

Exercises XII

January 20th

E12.1 *Mott scattering*

Consider an electron, confined to a volume V , in an external field $A_\mu(x)$. As a basis for the free solutions of the Dirac equation with momentum p and spin s we choose

$$\psi(x) = \sqrt{\frac{mc^2}{EV}} u(p, s) e^{-\frac{i}{\hbar}\langle p, x \rangle} \text{ (electron)}, \quad \psi(x) = \sqrt{\frac{mc^2}{EV}} v(p, s) e^{\frac{i}{\hbar}\langle p, x \rangle} \text{ (positron)}$$

with $(v(p, s)$ is $u(p, s)$, the first two components interchanged with the last two)

$$u(p, \frac{1}{2}) = \sqrt{\frac{E+mc^2}{2mc^2}} \begin{pmatrix} 1 \\ 0 \\ \frac{cp^z}{E+mc^2} \\ \frac{cp^+}{E+mc^2} \end{pmatrix}, \quad u(p, -\frac{1}{2}) = \sqrt{\frac{E+mc^2}{2mc^2}} \begin{pmatrix} 0 \\ 1 \\ \frac{cp^-}{E+mc^2} \\ \frac{-cp^z}{E+mc^2} \end{pmatrix}, \quad p^\pm = p^x \pm ip^y.$$

We consider the external field $A_\mu(x)$ only in lowest order perturbation theory (Born approximation). Then the transition matrix element between ψ_i and ψ_f is

$$S_{fi} = -\frac{ie}{\hbar c} \int d^4x \bar{\psi}_f(x) \gamma^\mu A_\mu(x) \psi_i(x).$$

- (1) Show that the matrix element is in principle determined by the Fourier transform of the four potential, taken at the four momentum transfer:

$$S_{fi} = -\frac{ie}{V} \frac{mc}{\hbar} \frac{1}{\sqrt{E_i E_f}} (\bar{u}_f \gamma^\mu u_i) (2\pi)^2 (\mathcal{F} A_\mu) \left(\frac{p_i - p_f}{\hbar} \right).$$

- (2) In particular, consider a Coulomb potential $A_0(x) = -\frac{Ze}{r}$, $\vec{A} = 0$. Using the momentum transfer $\vec{q} = \vec{p}_i - \vec{p}_f$, show that

$$S_{fi} = i4\pi \frac{Ze^2}{V} \frac{mc^2}{\sqrt{E_i E_f}} (\bar{u}_f \gamma^0 u_i) \frac{\hbar^2}{|\vec{q}|^2} 2\pi \delta(E_f - E_i).$$

- (3) Why does $N = \frac{V d^3 p_f}{(2\pi\hbar)^3}$ give the number of states in the momentum interval $d^3 p_f$? Thus, the transition probability per particle is given by $N|S_{fi}|^2$.

- (4) Compute the number R of transitions to the momentum interval $d^3 p_f$ per time:

$$R = \frac{4Z^2 \alpha^2 (\hbar c)^2}{V} \frac{m^2 c^4}{E_i E_f} |\bar{u}_f \gamma^0 u_i|^2 \frac{1}{|\vec{q}|^4} d^3 p_f \delta(E_f - E_i).$$

- (5) Why is $v_i = \frac{|\vec{p}_i|c^2}{E}$ the velocity of incoming particles and $j = \frac{v_i}{V}$ their current density? Compute the differential cross section

$$\frac{d\sigma}{d\Omega} = \frac{4Z^2 \alpha^2 m^2 (\hbar c)^2}{|\vec{q}|^4} |\bar{u}_f \gamma^0 u_i|^2.$$

- (6) In most experiments one will not know the spin polarizations. Thus we have to sum the outgoing states and to average the incoming ones:

$$\frac{d\sigma}{d\Omega} = \frac{4Z^2 \alpha^2 m^2 (\hbar c)^2}{|\vec{q}|^4} \frac{1}{2} \sum_{s_f, s_i} |\bar{u}_f \gamma^0 u_i|^2.$$

Show that this averaging can be written as a trace:

$$\frac{d\sigma}{d\Omega} = \frac{4Z^2\alpha^2 m^2 (\hbar c)^2}{2|\vec{q}|^4} \text{Tr} \left(\gamma^0 \frac{\gamma_\mu p_i^\mu + mc}{2mc} \gamma^0 \frac{\gamma_\mu p_f^\mu + mc}{2mc} \right).$$

- (7) Why does the trace of a product of an odd number of γ^μ vanish? Show the identity $\text{Tr}(a_\mu \gamma^\mu b_\nu \gamma^\nu) = 4\langle a, b \rangle$ and generalize it. Deduce

$$\frac{d\sigma}{d\Omega} = \frac{Z^2\alpha^2 \hbar^2}{2|\vec{q}|^4 c^2} (8E_i E_f - 4\langle p_i, p_f \rangle c^2 + 4m^2 c^4).$$

- (8) Show that $\langle p_i, p_f \rangle = \frac{E^2}{c^2} - p^2 \cos \vartheta$ and $|\vec{q}|^2 = 4p^2 \sin^2 \frac{\vartheta}{2}$, where $E_i = E_f = E$, $|\vec{p}_i| = |\vec{p}_f| = p$ and ϑ is the scattering angle. Deduce Mott's cross section ($\beta = \frac{v}{c}$):

$$\frac{d\sigma}{d\Omega} = \frac{Z^2\alpha^2 \hbar^2}{4p^2 \beta^2 \sin^4 \frac{\vartheta}{2}} \left(1 - \beta^2 \sin^2 \frac{\vartheta}{2} \right) = \left(\frac{d\sigma}{d\Omega} \right)_{\text{Rutherford}} \left(1 - \beta^2 \sin^2 \frac{\vartheta}{2} \right).$$

Homework XII

Return: January 27th

H12.1 Pair creation

We would like to consider a process similar to E12.1, namely the creation of an electron–positron pair by an external field.

- (1) Show that the matrix element is in principle determined by the Fourier transform of the four-potential, taken at the total four-momentum $P = p + p'$:

$$S_{fi} = -\frac{ie}{V} \frac{mc}{\hbar} \frac{1}{\sqrt{EE'}} \bar{v}(p', s') \gamma^\mu u(p, s) (2\pi)^2 (\mathcal{F}A_\mu) \left(\frac{P}{\hbar} \right).$$

The electron has four momentum p , spin s and spinor u , the positron p' , s' , v .

- (2) Perform the sum over the particle spins in $|S_{fi}|^2$ and show using $K = \frac{P}{\hbar}$:

$$|S_{fi}|^2 = \frac{(2\pi)^4 e^2 m^2 c^2}{\hbar^2 V^2 EE'} (\mathcal{F}A_\mu)(K) (\mathcal{F}A_\nu)(-K) \text{Tr} \left(\gamma^\mu \frac{\gamma^\kappa p_\kappa + mc}{2mc} \gamma^\nu \frac{\gamma^\kappa p'_\kappa - mc}{2mc} \right).$$

- (3) Compute the trace using the rules from E12.1.7 and show

$$|S_{fi}|^2 = \frac{(2\pi)^4 e^2}{\hbar^2 V^2 EE'} (\mathcal{F}A_\mu)(K) (\mathcal{F}A_\nu)(-K) (-m^2 c^2 g^{\mu\nu} + p^\mu p'^\nu + p^\nu p'^\mu - g^{\mu\nu} \langle p, p' \rangle).$$

- (4) Why is $N = \frac{V^2 d^3 p d^3 p'}{(2\pi\hbar)^6}$ the number of states in the momentum interval $d^3 p d^3 p'$?

- (5) Compute the Fourier transform for $A_0 = 0$, $\vec{A} = (a \cos \omega t) \vec{e}_z$, $a \in \mathbb{R}$.

- (6) Calculate the number of produced pairs as $N|S_{fi}|^2$:

$$\frac{e^2}{24\pi\hbar^2 c^3} |a|^2 \omega^2 \sqrt{1 - \frac{4m^2 c^4}{\hbar^2 \omega^2}} \left(1 + \frac{2m^2 c^4}{\hbar^2 \omega^2} \right) (2\pi)^4 \delta \left(\frac{\vec{p} + \vec{p}'}{\hbar} \right) \delta \left(\frac{E + E'}{\hbar} - \omega \right).$$

- (7) Argue why the number R of pairs per volume and time unit is

$$R = \frac{e^2}{24\pi\hbar^2 c^3} |a|^2 \omega^2 \sqrt{1 - \frac{4m^2 c^4}{\hbar^2 \omega^2}} \left(1 + \frac{2m^2 c^4}{\hbar^2 \omega^2} \right).$$

- (8) How does the angular distribution of the particle momenta look like?

- (9) In which cases can pair production take place at all?

(30 points)