

D-Branes in a Non-Compact WZNW Model

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Motivation . . .

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» . . . and Outline

The H_3^+ Model

Cardy-Lewellen Constraints

Implementation of Factorization
Constraint in H_3^+

The ' $b^{-2}/2$ '-Constraint

Epilogue

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- String Theory in gravitational backgrounds needs to be understood.

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- String Theory in gravitational backgrounds needs to be understood.
- 3D (Euclidean) Anti-de-Sitter: Feasible Background

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- String Theory in gravitational backgrounds needs to be understood.
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- Also: AdS_3/CFT_2 Correspondence

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- Focus on corresponding CFT (H_3^+ Model): Interesting in its own right!

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- Focus on corresponding CFT (H_3^+ Model): Interesting in its own right!
- Reason: It is a *non-rational* CFT

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- Study couplings of closed strings to D-branes

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- However, in the (non-rational) H_3^+ CFT: Problems in retaining power of factorization constraint.

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- We want to discuss today, what problems need to be overcome and ...

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- This works very well in *rational* CFT
- However, in the (non-rational) H_3^+ CFT: Problems in retaining power of factorization constraint.
- We want to discuss today, what problems need to be overcome and ...
- ... how power of factorization constraint is maintained.

The H_3^+ Model

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- $\mathbf{H}_3^+ = \text{Euclidean } AdS_3:$

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- $H_3^+ = \text{Euclidean } AdS_3:$
- $AdS_3 : X_0^2 - X_1^2 - X_2^2 + X_3^2 = 1$

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■ $H_3^+ =$ Euclidean AdS_3 :

■ $AdS_3 : X_0^2 - X_1^2 - X_2^2 + X_3^2 = 1$

$$\begin{pmatrix} X_0 + X_1 & X_2 + X_3 \\ X_2 - X_3 & X_0 - X_1 \end{pmatrix} \in SL(2, \mathbb{R})$$

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■ Thus, $H_3^+ = \{h \in SL(2, \mathbb{C}) : h = h^\dagger, \text{tr}(h) > 0\}$.

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- Thus, $H_3^+ = \{h \in SL(2, \mathbb{C}) : h = h^\dagger, \text{tr}(h) > 0\}$.
- Consequently, an action of $SL(2, \mathbb{C})$ is admitted:

$$h \in H_3^+ \mapsto ghg^\dagger \in H_3^+ \text{ for } g \in SL(2, \mathbb{C})$$

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- Central Theme: Making use of $SL(2, \mathbb{C})$ symmetry.

The H_3^+ Bulk CFT

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- Now, consider closed strings moving in H_3^+ :

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	flat space - free boson CFT	euclid. $AdS_3 - H_3^+$ CF
Equation of Motion		
Symmetry		
Mode Expansion		
Representations		
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Sugawara Form	$T(z) = \frac{1}{2} : (\partial\phi)(z) (\partial\phi)(z) :$	$T(z) = \frac{b^2}{2} : J^a(z) J^a(z) :$

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- Important piece of structure data: Bulk 3-point function coefficient

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$$C(j_3, j_2, j_1) \propto \langle \Theta_{j_3}(u_3|z_3) \Theta_{j_2}(u_2|z_2) \Theta_{j_1}(u_1|z_1) \rangle$$

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- For doubling trick to work, we need a *gluing condition*:

$$J^a(z) = \rho^a_b \bar{J}^b(\bar{z}) \quad \text{such that} \quad T(z) = \bar{T}(\bar{z}) \quad \text{at} \quad z = \bar{z}$$

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- \mathbf{H}_3^+ : Two choices for ρ . We consider $\rho \bar{J}^3 = \bar{J}^3$, $\rho \bar{J}^\pm = -\bar{J}^\pm$.

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 \Rightarrow so-called AdS_2 branes; labelled by a parameter $\alpha \in \mathbb{R}$.

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$$J^a(z) = \rho^a_b \bar{J}^b(\bar{z}) \quad \text{such that} \quad T(z) = \bar{T}(\bar{z}) \quad \text{at} \quad z = \bar{z}$$

- \mathbf{H}_3^+ : Two choices for ρ . We consider $\rho \bar{J}^3 = \bar{J}^3$, $\rho \bar{J}^\pm = -\bar{J}^\pm$.
 \Rightarrow so-called AdS_2 branes; labelled by a parameter $\alpha \in \mathbb{R}$.
- 1-point correlator:

$$\langle \Theta_j(u|z) \rangle_\alpha \equiv \langle \theta_j(u|z) \bar{\theta}_j(\bar{u}|\bar{z}) \rangle_\alpha = |z - z^*|^{-2h(j)} |u + u^*|^{2j} A_\sigma(j|\alpha)$$

Boundary CFT

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$$\langle \Theta_j(u|z) \Psi_{j'}(t|x) \rangle_\alpha \propto C(j, j'|\alpha)$$

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- For $j' = 0$, i.e. $\Psi_{j'=0} = \mathbb{1}$: $C(j, 0|\alpha) = A(j|\alpha)$



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Factorization Constraint

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$$\sum_q C(2, 1, q) F_{21, \bar{2}\bar{1}}^q(z) A(q|\alpha) = \sum_p C(2, p|\alpha) F_{2\bar{2}, 1\bar{1}}^p(1-z) C(1, p|\alpha)$$

with crossing ratio $z = \frac{|z_2 - z_1|^2}{|z_2 - z_1^*|^2}$

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- ... the sum is **finite**, ...
- ... $C(i, 0|\alpha) = A(i|\alpha)$ is always **well-defined**, and ...
- ... on the LHS, the **limit** $z \rightarrow 1$ can be taken.

Implementation of Factorization Constraint in H_3^+

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- H_3^+ is non-rational.

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Epilogue

- H_3^+ is **non-rational**.
- Between two generic H_3^+ field operators: **Continuous** spectrum of propagating fields

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Epilogue

- H_3^+ is **non-rational**.
- Between two generic H_3^+ field operators: **Continuous** spectrum of propagating fields
- The **identity** is not guaranteed to propagate in open string channel.

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- Nevertheless, assuming a **discrete** open string spectrum, $C(j_s, 0|\alpha) = A(j_s|\alpha)$ remains valid and well-defined.
- \Rightarrow **Need to distinguish two different kinds of D-branes: Continuous and Discrete.**

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Epilogue

- Using singular field $\Theta_{1/2}$ and generic Θ_j , i.e. consider

$$G_{j,\alpha}^{(2)}(u_i|z_i) := \langle \Theta_{1/2}(u_2|z_2)\Theta_j(u_1|z_1) \rangle_\alpha$$

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- The Sugawara Construction implies

$$L_{-1} = -b^2 \eta_{ab} J_{-1}^a J_0^b$$

The '1/2'-Constraint (1)

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- The Sugawara Construction implies the Knizhnik-Zamolodchikov equation

$$\left\langle \left\{ \left(L_{-1} + b^2 \eta_{ab} J_{-1}^a J_0^b \right) \Theta_{1/2}(u_2|z_2) \right\} \Theta_j(u_1|z_1) \right\rangle_\alpha = 0$$

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- $C_\pm(j) = C(j, \frac{1}{2}, j \pm \frac{1}{2})$

The '1/2'-Constraint (2)

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The '1/2'-Constraint (2)

- The conformal blocks $\mathcal{F}_{j,+}^s(u|z)$, $\mathcal{F}_{j,-}^s(u|z)$ depend on the crossing ratios

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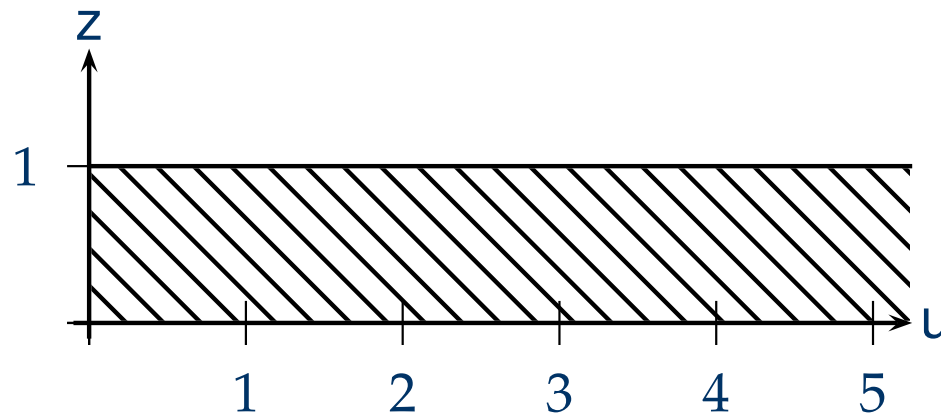
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- Limit $z \rightarrow 1$ can be taken without problems

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Epilogue

- Limit $z \rightarrow 1$ can be taken without problems
- Assuming $C(2, 1|\alpha) = A(2|\alpha)$, i.e. discrete open string spectrum

$$\Rightarrow A_{\text{disc.}}(j|m, n) \quad \text{with} \quad m \in \mathbb{Z}$$

$$\Rightarrow j_{m,n} = -\frac{1}{2} + \frac{m}{2} + b^{-2} \frac{n}{2}$$

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- Compare to the rational case (Cardy):
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- Assuming continuous open string spectrum

$$\Rightarrow A_{\text{cont.}}(j|\alpha) \quad \text{with} \quad \alpha \in \mathbb{R}$$

$$\Rightarrow j_\alpha = -\frac{1}{2} + i\alpha$$

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- ... the solutions $A_{\text{disc.}}(j|m, n)$ and $A_{\text{cont.}}(j|\alpha)$ are not unique!

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- But ...
- ... the solutions $A_{\text{disc.}}(j|m, n)$ and $A_{\text{cont.}}(j|\alpha)$ are not unique!
- Need second constraint to fix leftover arbitrariness.

The ' $b^{-2}/2$ '-Constraint

Knizhnik-Zamolodchikov equation

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Epilogue

■ Consider

$$\langle \Theta_{b^{-2}/2}(u_2|z_2) \Theta_j(u_1|z_1) \rangle_\alpha$$

Knizhnik-Zamolodchikov equation

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- Again Knizhnik-Zamolodchikov equation. It is solved by

$$G_{j,\alpha}^{(2)}(u_i|z_i) \propto \sum_{\epsilon=+,-,\times} C_\epsilon(j) A_\sigma(j_\epsilon|\alpha) \mathcal{F}_{j,\epsilon}^s(u|z)$$

Knizhnik-Zamolodchikov equation

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$$G_{j,\alpha}^{(2)}(u_i|z_i) \propto \sum_{\epsilon=+,-,\times} C_\epsilon(j) A_\sigma(j_\epsilon|\alpha) \mathcal{F}_{j,\epsilon}^s(u|z)$$

with $C_\epsilon(j) = C(j, \frac{1}{2}, j_\epsilon)$, where

$$j_\pm = j \pm \frac{b^{-2}}{2} \quad \text{and} \quad j_\times = -j - 1 - \frac{b^{-2}}{2}$$

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Epilogue

■ Crossing Ratios

$$z := \frac{|z_2 - z_1|^2}{|z_2 - \bar{z}_1|^2} \quad \text{and} \quad u := \frac{|u_2 - u_1|^2}{|u_2 + \bar{u}_1|^2}$$

$$■ \quad 0 \leq z \leq 1 \quad \text{and} \quad 0 \leq u < \infty$$

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- $0 \leq z \leq 1$ and $0 \leq u < \infty$
- $\mathcal{F}_{j,+}^s(u|z), \mathcal{F}_{j,-}^s(u|z), \mathcal{F}_{j,\times}^s(u|z)$ are simultaneously well-defined only for $z < u < 1$.

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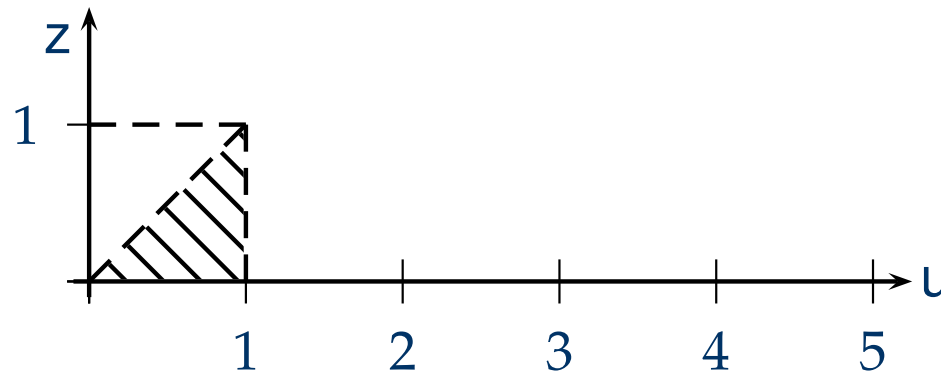
» Limit $z \rightarrow 1$

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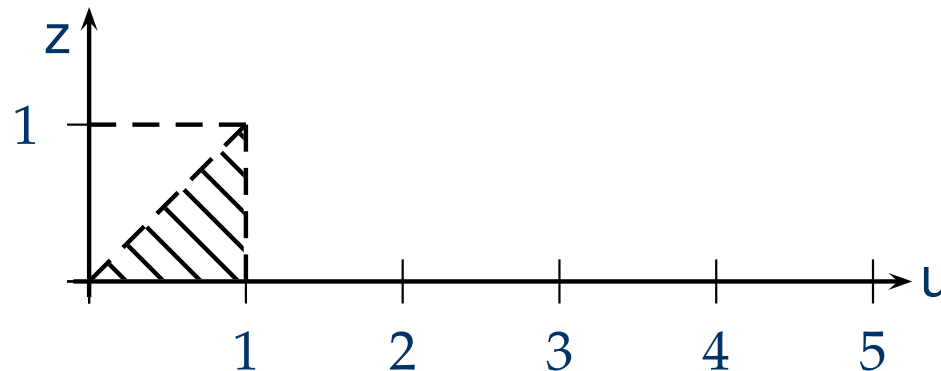
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- $0 \leq z \leq 1$ and $0 \leq u < \infty$
- $\mathcal{F}_{j,+}^s(u|z), \mathcal{F}_{j,-}^s(u|z), \mathcal{F}_{j,\times}^s(u|z)$ are simultaneously well-defined only for $z < u < 1$.



- Thus, limit $z \rightarrow 1$ cannot be taken! No constraint derivable!

Resolution

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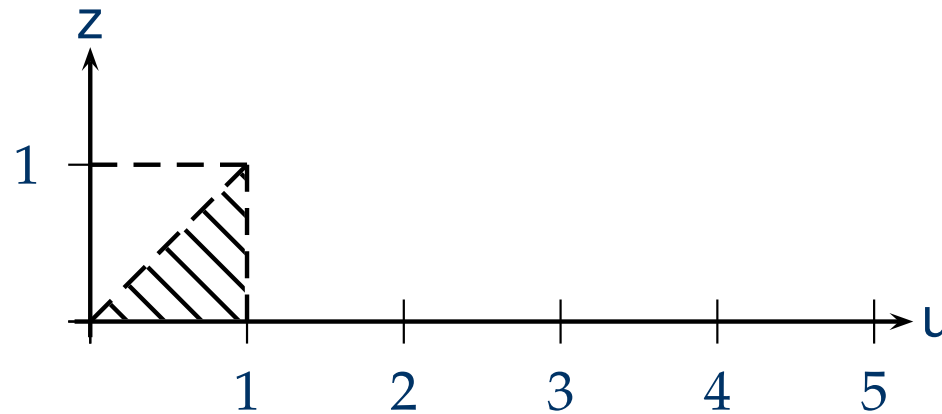
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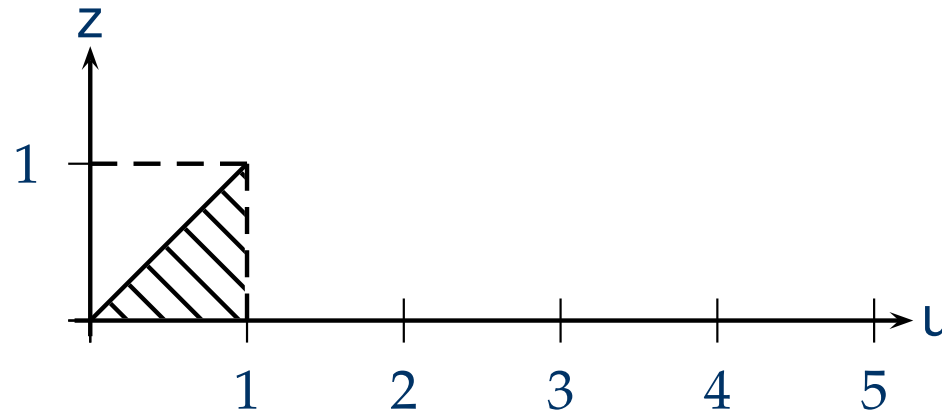
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Epilogue



- Our solution for $G_{j,\alpha}^{(2)}(u_i|z_i)$ is finite at $u = z$.

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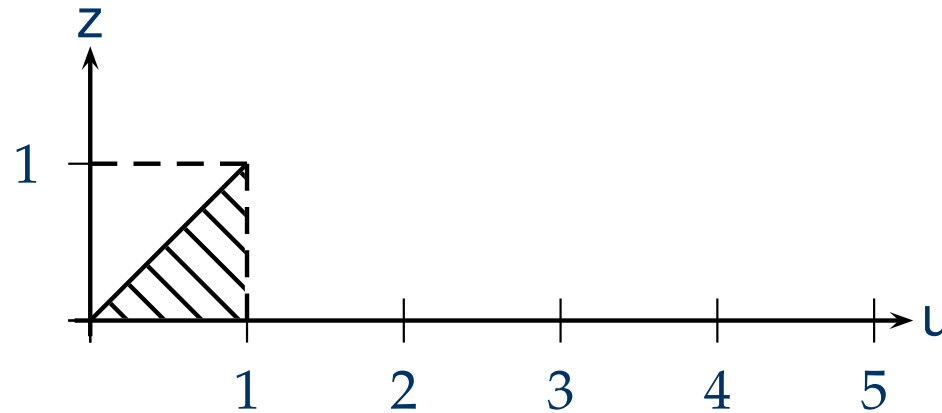
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Epilogue



- Our solution for $G_{j,\alpha}^{(2)}(u_i|z_i)$ is finite at $u = z$.
- But we observe that $\mathcal{F}_{j,-}^s(u|z)$ and $\mathcal{F}_{j,\times}^s(u|z)$ become linearly dependent at $u = z$.

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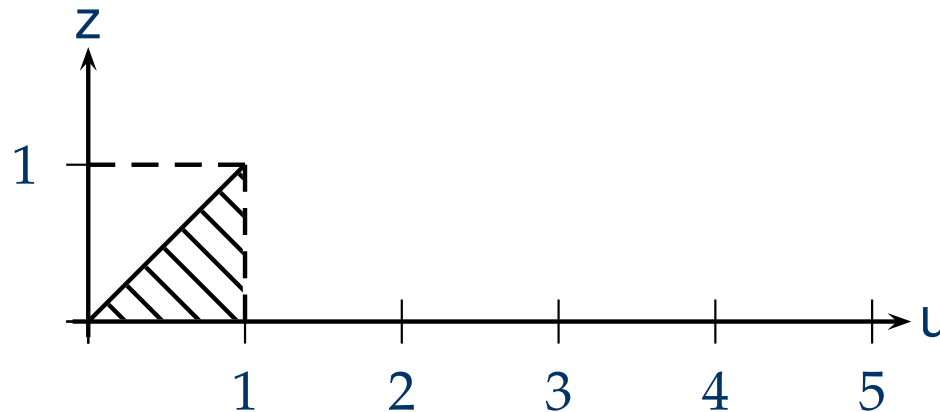
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- But we observe that $\mathcal{F}_{j,-}^s(u|z)$ and $\mathcal{F}_{j,\times}^s(u|z)$ become linearly dependent at $u = z$.
- $u = z$ suggested to be a "new type of singularity" (*Hosomichi-Ribault proposal*)

Solution in Second Patch

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Epilogue

- We can also give a basis-set of solutions $\tilde{\mathcal{F}}_{j,+}^s(u|z)$, $\tilde{\mathcal{F}}_{j,-}^s(u|z)$ and $\tilde{\mathcal{F}}_{j,\times}^s(u|z)$ which are well-defined for $u < z < 1$.

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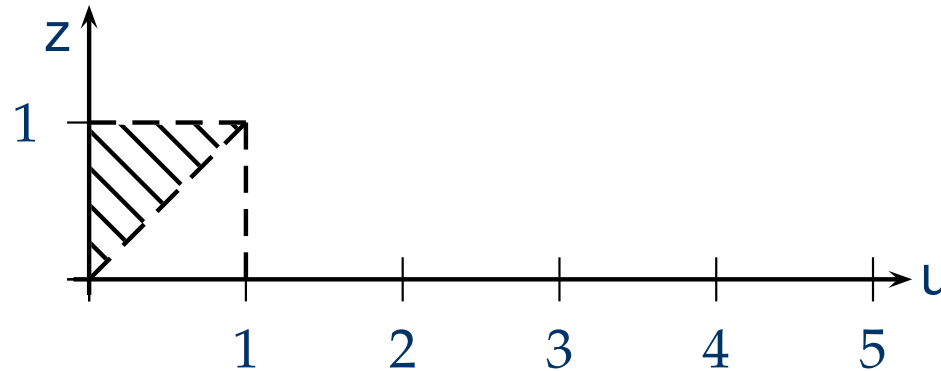
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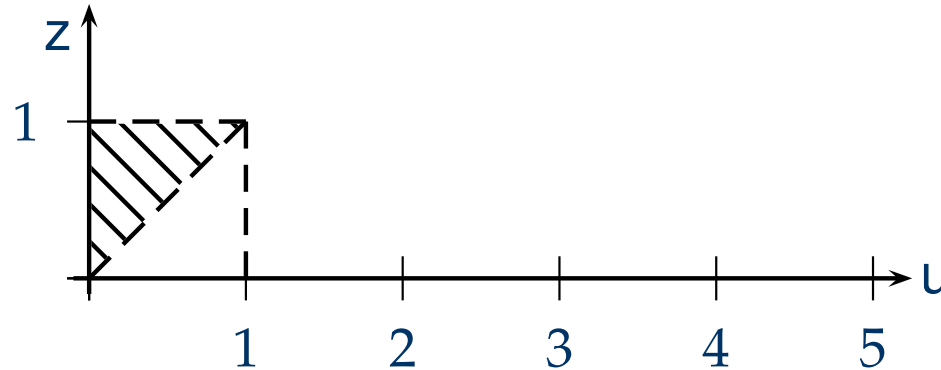
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- Solution in region $u < z < 1$ is therefore expanded as

$$\tilde{G}_{j,\alpha}^{(2)}(u_i|z_i) \propto \sum_{\epsilon=+,-,\times} \tilde{a}_{\epsilon}^j(\alpha) \tilde{\mathcal{F}}_{j,\epsilon}^s(u|z)$$

Solution in Second Patch

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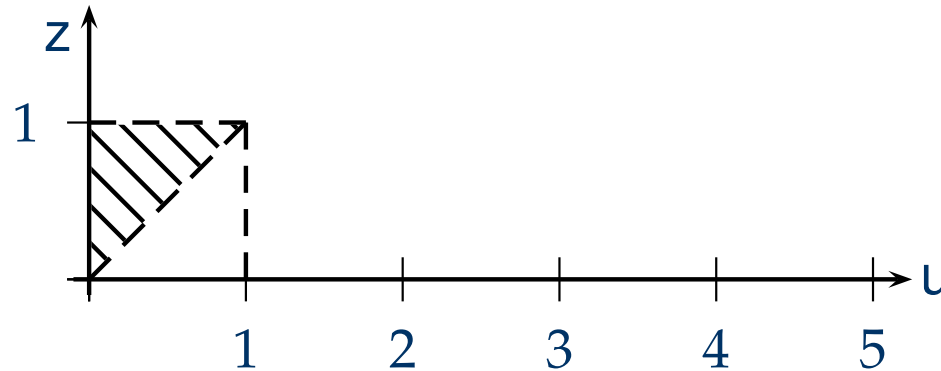
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Epilogue

- We can also give a basis-set of solutions $\tilde{\mathcal{F}}_{j,+}^s(u|z)$, $\tilde{\mathcal{F}}_{j,-}^s(u|z)$ and $\tilde{\mathcal{F}}_{j,\times}^s(u|z)$ which are well-defined for $u < z < 1$.



- Solution in region $u < z < 1$ is therefore expanded as

$$\tilde{G}_{j,\alpha}^{(2)}(u_i|z_i) \propto \sum_{\epsilon=+,-,\times} \tilde{a}_{\epsilon}^j(\alpha) \tilde{\mathcal{F}}_{j,\epsilon}^s(u|z)$$

- Again, it is finite at $u = z$.

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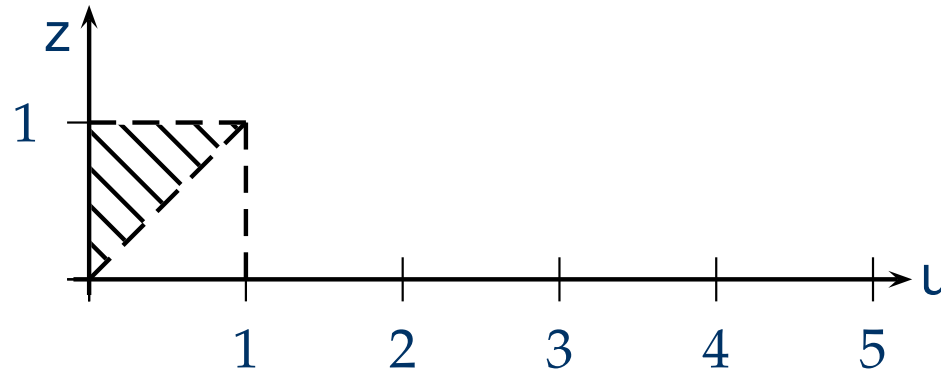
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$$\tilde{G}_{j,\alpha}^{(2)}(u_i|z_i) \propto \sum_{\epsilon=+,-,\times} \tilde{a}_\epsilon^j(\alpha) \tilde{\mathcal{F}}_{j,\epsilon}^s(u|z)$$

- Again, it is finite at $u = z$.
- But also again, $\tilde{\mathcal{F}}_{j,-}^s(u|z)$ and $\tilde{\mathcal{F}}_{j,\times}^s(u|z)$ become linearly dependent at $u = z$.

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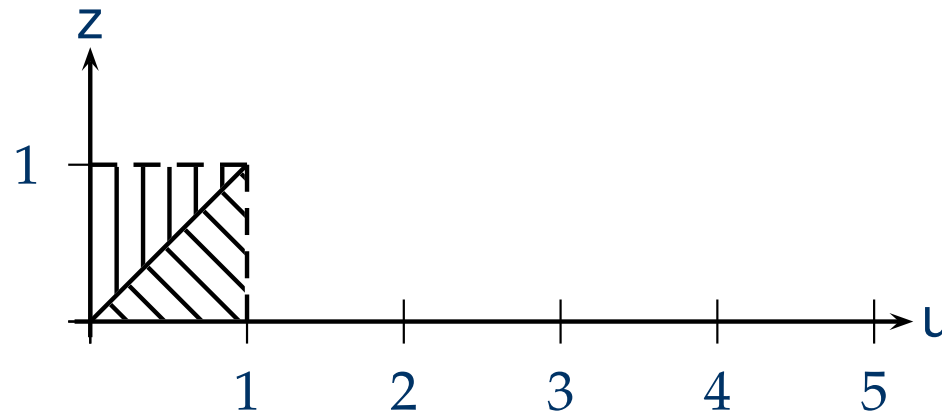
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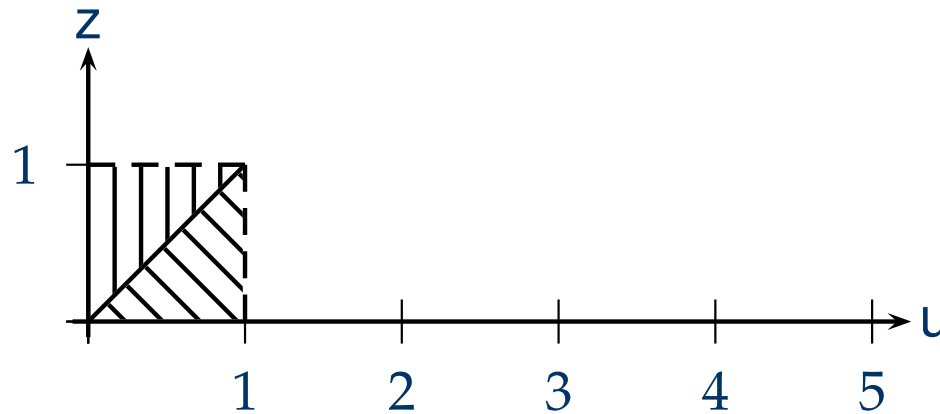
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- We show that the two solutions can be matched at $u = z$.

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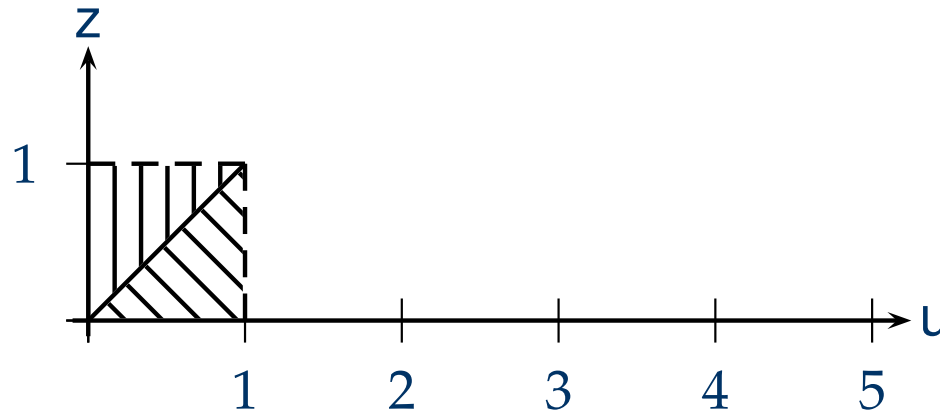
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- We show that the two solutions can be matched at $u = z$.
- From this, only the coefficient $\tilde{a}_+^j(\alpha)$ can be uniquely fixed ...

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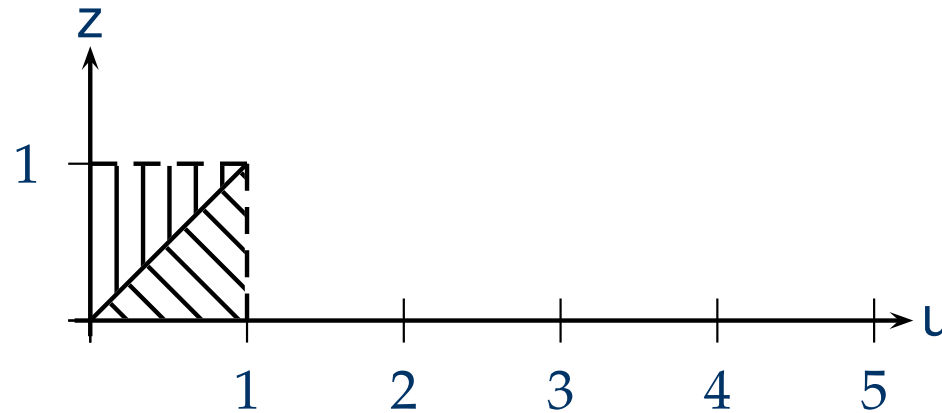
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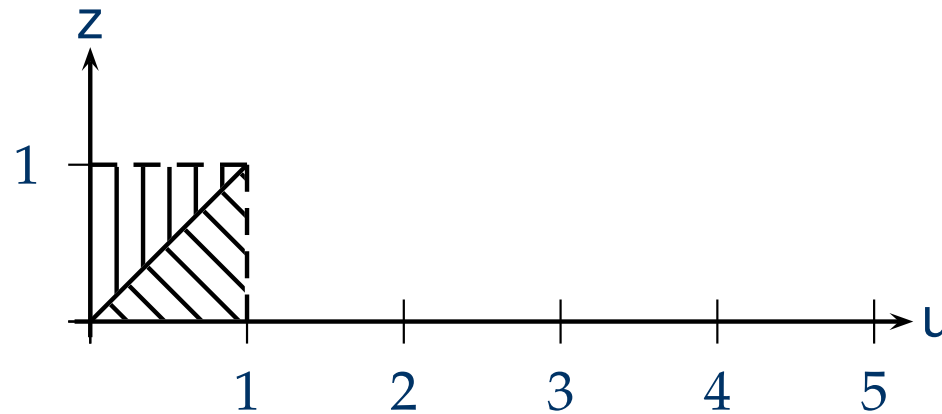
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- Thus: Obtain a one parameter set of possible two point functions in the patch $u < z < 1$.

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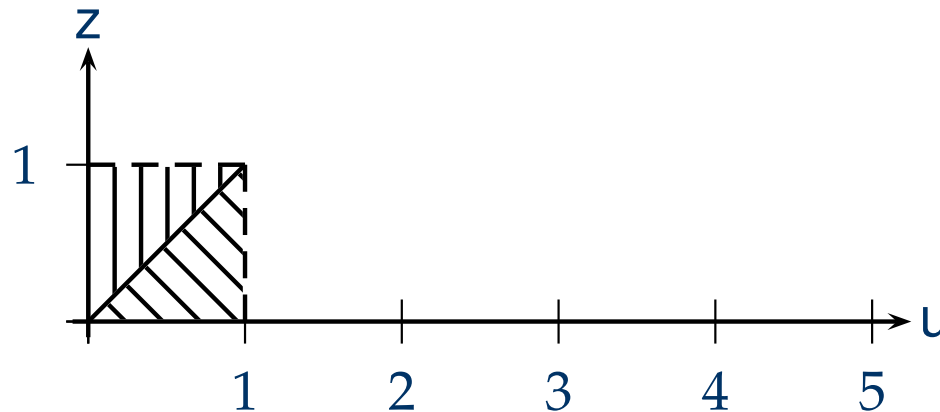
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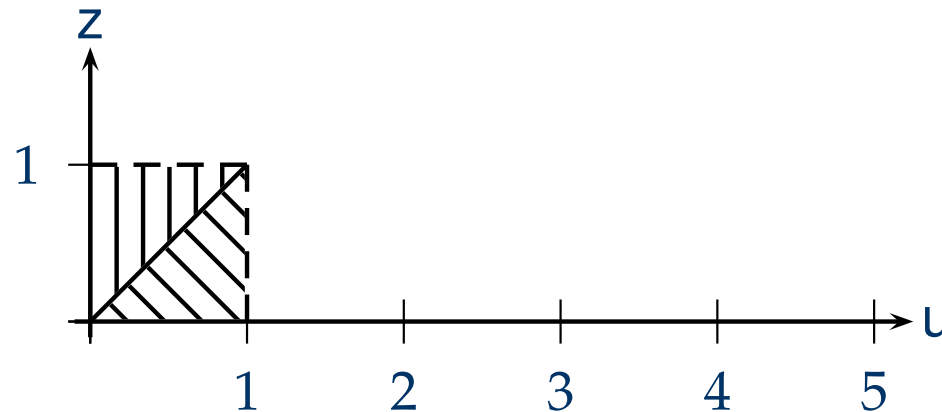
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- From this, only the coefficient $\tilde{a}_+^j(\alpha)$ can be uniquely fixed ...
- ... together with a linear combination of $\tilde{a}_-^j(\alpha)$ and $\tilde{a}_\times^j(\alpha)$.
- Thus: Obtain a one parameter set of possible two point functions in the patch $u < z < 1$.
- This means, the limit $z \rightarrow 1$ may not be fixed uniquely.
- Henceforth, it seems that the factorization constraint loses its power.

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BUT ...

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BUT ...

- ... in the limit $z \rightarrow 1$, the two blocks $\tilde{\mathcal{F}}_{j,-}^s(u|z)$ and $\tilde{\mathcal{F}}_{j,\times}^s(u|z)$ do not contribute to the identity channel!

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- ... in the limit $z \rightarrow 1$, the two blocks $\tilde{\mathcal{F}}_{j,-}^s(u|z)$ and $\tilde{\mathcal{F}}_{j,\times}^s(u|z)$ do not contribute to the identity channel!
- Only the uniquely fixed part involving $\tilde{\mathcal{F}}_{j,+}^s$ does!

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- Arbitrariness vanishes.

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- Arbitrariness vanishes.
- Factorization constraint remains powerful and predictive!

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- Only the uniquely fixed part involving $\tilde{\mathcal{F}}_{j,+}^s$ does!
- Arbitrariness vanishes.
- Factorization constraint remains powerful and predictive!
- New shift equations are in the discrete case:

$$A_\sigma(j|\alpha)A_\sigma\left(\frac{1}{2b^2}|\alpha\right) = \frac{1 + (2j + 1)b^2}{1 + b^2}A_\sigma\left(j + \frac{1}{2b^2}|\alpha\right)$$

and in the continuous case:

$$\text{Res}_{j_s=b^{-2}/2}C(j_s, 0|\alpha)A_\sigma(j|\alpha) = (1 + (2j + 1)b^2)A_\sigma\left(j + \frac{1}{2b^2}|\alpha\right)$$



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Summary

- We have discussed the factorization constraint in the H_3^+ model, a specific non-rational CFT.

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- Compared to rational CFT, the factorization constraint is not straightforwardly implemented.

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- We have discussed the factorization constraint in the H_3^+ model, a specific non-rational CFT.
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- Need to distinguish two types of D-branes: Continuous and Discrete.
- A new kind of "singularity" ($u = z$ singularity) appears when considering the $b^{-2}/2$ -constraint.
- Its existence has been suggested by a mapping to Liouville theory (Hosomichi-Ribault proposal).

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Our Results

- We can show the emergence of the $u = z$ singularity in the H_3^+ model by an explicit calculation.

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- We can show the emergence of the $u = z$ singularity in the H_3^+ model by an explicit calculation.
- Treating the singularity as suggested by the work of Hosomichi and Ribault, we obtain a new explicit factorization constraint.

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- We show that the 1-point amplitudes for the discrete D-branes $A_{\text{disc.}}(j|m, n)$ with $m \in \mathbb{Z}$ are further constrained to also having $n \in \mathbb{Z}$.

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brane spectrum \longleftrightarrow representations of symmetry algebra

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- We show that the 1-point amplitudes for the discrete D-branes $A_{\text{disc.}}(j|m, n)$ with $m \in \mathbb{Z}$ are further constrained to also having $n \in \mathbb{Z}$.
- This fits in well with the Cardy interpretation brane spectrum \longleftrightarrow representations of symmetry algebra
- We show that the 1-point amplitudes for the continuous 1-point functions $A_{\text{cont.}}(j|\alpha)$ with $\alpha \in \mathbb{R}$ obey new constraint without any new restrictions.

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- Why did the factorization constraint not lose its power?

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- Why did the factorization constraint not lose its power?
- Can you see that in Liouville Theory?

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- Why did the factorization constraint not lose its power?
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- "Singularity" at $u = z$ specific to H_3^+ Model. What generic features remain in more general non-rational CFTs?

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- Besides the AdS_2 branes, there is a second class: S^2 branes.

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- Besides the AdS_2 branes, there is a second class: S^2 branes.
- What happens when studying the $b^{-2}/2$ constraint for them? (Mapping to Liouville Theory might not be possible...)

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